Cosmic-ray Hadrons at GeV—TeV energies and CALET

Nicholas Cannady ISCRA 2022 Erice, Italy Some species in the CR flux are overrepresented compared to local Galactic (i.e. Solar System) abundances



Cycle of matter in the Galaxy



Abundances of cosmic-ray species



Abundances as measured at Earth

Abundances of cosmic-ray species



Abundances as measured at Earth



Normalized to Solar System abundances

Abundances of cosmic-ray species



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OB Association model of injection



Synthesis of MSM in OB Associations

OB Association model of injection



Synthesis of MSM in OB Associations Condensation onto grains, along with other nuclei and p

Collection of charge on outside of grains -> increase rigidity for same momentum



e.g. Ivlev et al. 2018

OB Association model of injection



Synthesis of MSM in OB Associations

Charge sputtering cross section Lingenfelter et al. 2019

Condensation onto grains, along with other nuclei and p

Collection of charge on outside of grains -> increase rigidity for same momentum



e.g. lvlev et al. 2018



Efficient acceleration of grains to supra-thermal energies

Grains are dissociated through collisions, fragments accelerated to CR energies



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Stages of GCR lifetime



Synthesis:

Stellar fusion processes, massive star ejecta, explosive nucleosynthesis

Injection:

Dust grain acceleration to supra-thermal energies in strong shocks

Acceleration:

Grain sputtering and acceleration to GCR energies in strong shocks **Propagation:**

Accelerated GCRs escape local environment and transport throughout the Galaxy Some species in the CR flux are overrepresented compared to local Galactic (i.e. Solar System) abundances



Not due to enrichment of MSM, must arise from fragmentation of heavier nuclei producing significant secondary component



For a nuclear species i,

$$\frac{dN_i}{dt} = D\nabla^2 N_i + \frac{\partial}{\partial E} [bN_i] + Q_i - \frac{N_i}{\tau_i} + \sum_{j>i} \frac{P_{ji}}{\tau_j} N_j$$

Let's look at the effect of primary spallation into secondaries: Neglect diffusion, energy losses, and source injection

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For a primary species P,	$\frac{dN_P}{dt} = -$	$-rac{N_P}{ au_P}$	Assumption: nothing significantly contributing to secondary P
For a secondary species S,	$\frac{dN_S}{dt} = -$	$-\frac{N_S}{\tau_S} + \frac{P_{PS}}{\tau_P} N_P$	

For a **primary** species P, $\frac{dN_P}{dt} = -\frac{N_P}{\tau_P}$ Assumption: nothing significantly contributing to secondary P For a **secondary** species S, $\frac{dN_S}{dt} = -\frac{N_S}{\tau_S} + \frac{P_{PS}}{\tau_P} N_P$ Change independent variable: $\xi = \rho v$ $\downarrow t \rightarrow \xi$ $\tau_P \rightarrow \xi_P$ $\tau_S \rightarrow \xi_S$

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$$\frac{dN_P}{d\xi} = -\frac{N_P}{\xi_P} \qquad \qquad N_P(\xi) = N_{P0}e^{-\xi/\xi_P} \qquad (N_P(0) = N_{P0})$$

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For a **primary**

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Change independent variable: $\xi = \rho v$

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t

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$$\frac{N_S}{N_P} = \frac{P_{PS}\xi_S}{\xi_S - \xi_P} \left[\exp\left(\frac{\xi}{\xi_P} - \frac{\xi}{\xi_S}\right) - 1 \right]$$

Secondary-to-primary ratio for light nuclei



Secondary-to-primary ratio for light nuclei



More realistic results require (at least) reintroduction of the diffusion term to account for distribution of path lengths traveled by particles (rather than assuming the average)

Secondary-to-primary ratios

Beyond this small illustrative example, the real situation is much more complex:

- Diffusion, energy loss, and source terms that we neglected
 - Treatment of diffusion leads to different distributions of pathlengths which modify the resulting secondary/primary ratios
 - Isotropic diffusion Gaussian distribution of pathlengths
 - Leaky box model exponential distribution of pathlengths
- Energy dependence of parameters in the transport equation
 - Diffusion coefficient
 - Energy loss and source injection terms
- Varying levels of "primary-ness" for each species measured separately
 - Mostly pure primaries, e.g. p, He, C, O, ..., Si, ..., Fe
 - Mixed, e.g. N with significant primary component and secondary component from interactions of primary O
 - Mostly pure secondaries, e.g. Li, Be, B

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Energy-dependent measurement of elemental CR spectra up to the limits of Galactic accelerators are critical input for models of acceleration and propagation of GCRs

Cosmic-ray all-particle spectrum

GeV PeV TeV s GeV)-1 **Relatively featureless except for small** Fluxes of Cosmic Roys changes in spectral index lux (m' sr (1 particle per m²-second) 10 Direct measurements are limited to 10 ~several PeV total energies simply by feasibility of required detector size 10-10 Knee (1 particle per m²-year) 10 10 Only up to the knee, where the transition to extragalactic sources dominating the flux begins Ankle 10-25 (1 particle per km²-year) 10 1018 1018 10¹¹ 10¹² 10¹³ 10¹⁴ 10¹⁵ 10¹⁶ 10¹⁷

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CALET and DAMPE Calorimeters

CALET



DAMPE

Depth(EM) = 30 X_0 Depth(HAD) = 1.2 λ_p Geom. Fact. = 0.12 m² sr Depth(EM) = 32 X_0 Depth(HAD) = 1.6 λ_p Geom. Fact. = 0.3 m² sr

8/4/2022

PAMELA and AMS-02 Spectrometers



AMS-02

PAMELA

295 cm

ATIC and CREAM balloon-borne calorimeters



Instrument	Energy measurement technique	Charge range and resolution	Flight duration	Atmospheric depth ^a (g/cm ²)	Effective exposure (m ² -sr-days)	Observed number of protons >6 TeV
ATIC	Calorimeter (0.75 λ_I , 18 X_0)	$1 \leqslant Z \leqslant 28 \Delta Z = 0.3$	\sim 48 days	4.3 (7.9)	5	~720
TRACER ^a	TRD	$8 \leq Z \leq 28 \Delta Z = 0.3 (O) 0.5 (Fe)$ $3 < Z \leq 28 \Delta Z = 0.3 (O) 0.5 (Fe)$	~10 days ~4 days	3.9 (9.2)	50 20	None
CREAM	Calorimeter (0.5 λ_l , 20 X ₀)	$1 \leqslant Z \leqslant 28 \Delta Z = 0.2$	$\sim \! 160 \text{ days}$	3.9 (6.8)	48	>5000
	TRD	$3 < Z \le 28 \Delta Z = 0.2$	~42 days	3.9(7.9)	55	None
JACEE	Emulsion (\sim 0.05 λ_I , \sim 4 X ₀)	$1 \leqslant Z \leqslant 28$ Charge group	\sim 60 days (1436 m ² hr)	5.3 (28)	$\sim 10 \ (644 \ m^2 \ hr)$	656
RUNJOB	Emulsion ($\sim 0.2 \lambda_l$, $\sim 4 X_0$)	$1 \leqslant Z \leqslant 28$ Charge group	${\sim}60$ days (575 m ² hr)	10 (48)	6 (p); 24 (>C)	Close to JACEE

ISS-CREAM

Space-borne successor to balloonborne CREAM family of missions

- Silicon charge detector
- Carbon targets induce nuclear interactions
- Top and bottom counting detectors for e/p separation
- Scintillating fiber tungsten sheet sampling calorimeter for energy measurement
- Boronated scintillator detector for more e/p separation

Began operations August 2017 Stopped February 2019



The NUCLEON apparatus



- Charge measurement system four planes of pad silicon detectors (1.5×1.5 cm²) (1);
- tracker for KLEM energy measurement carbon target of 0.25 3 proton interaction lengths (2) and six planes of microstrip silicon detectors (0.4mm pitch) with tungsten between them (~2mm each, ~3 X-lengths in total) (3);
- trigger sysytem three double scintillator planes (4).

Active area 500*500 mm². Geometrical factor ~0.2 m²sr.

Ionization calorimeter (IC) (5) – six planes of tungsten absorber (~8mm each, ~12 X-lengths in total) with silicon strip detectors (1mm pitch).

Active area 250*250mm².

Geometrical factor (together with charge and KLEM systems) ~0.06 m²sr. ~0.2 m² sr for nuclei



10604 independent electronic channels in total

Hadrons in the CALET calorimeter

Flight data He candidate with 400 GeV total deposit in instrument



Nuclear interaction: $A + X \rightarrow$

- p, n, A'
- K^{\pm} , K^{0} , $\pi^{\pm} \rightarrow \mu^{\pm}$

• $\pi^0 \rightarrow \gamma \gamma$

Hadron analysis in CALET

- Careful validation of selection parameters is performed based on identical analysis in multiple simulation packages:
- EPICS/COSMOS
- Fluka
- Geant4

Incomplete list of cuts used:

- Charge selection in CHD
- Combinatorial Kalman Filter tracking
- Require **low** IMC concentration ~
- Consistency of TASC energy deposits with IMC reconstructed track
- TASC shower topology cuts to remove mis-reconstructed tracks

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CHD-Y Charge

EPICS

/GeV < 2000

Entries

104

0

10 15 20

25 30

Energy spectrum unfolding

Due to penetrating nature of many secondaries from nuclear interactions, much of the shower energy is carried out of CALET

Primary energy reconstruction is not possible with high precision, but energy spectra can be *unfolded* using iterative Bayesian procedure



Proton spectrum measurement



Protons

- Spectral hardening first reported by PAMELA
- CREAM-III reported hard spectrum at TeV energies
- AMS-02 confirmed spectral hardening
- CALET measures over full soft-to-hard transition
- DAMPE extends this to see clear softening
- Preliminary CALET updates show softening as well

Helium

- Spectral hardening also reported by PAMELA
- CREAM-I reported less prominent hardening than protons at TeV energies
- AMS-02 confirmed spectral hardening
- DAMPE measurement demonstrates spectral shape similar to protons
- CALET preliminary results confirm this

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Nuclear spectra through Fe



Carbon and Oxygen and C/O ratio



Carbon and Oxygen and C/O ratio



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Preliminary Boron-to-Carbon ratio



Iron and Nickel and Ni/Fe ratio



Primary and secondary spectral shapes with AMS-02



Primary and secondary spectral shapes with AMS-02

Aguilar et al. 2021 PRL 126, 041104



What's coming next?



Summary

- Element-resolved measurements of GCRs support OB association origins
- \rightarrow Separation of elements based on T_c when normalized 81% SS + 19% MSM
 - ightarrow Dividing temperature ~1250K consistent with OB associations
 - \rightarrow Enrichment abundances of MSM from Wolf-Rayet winds and SNe
- \rightarrow Grain acceleration as injection mechanism favored by refractory enhancement
 - \rightarrow Pattern across elements follows Z^{2/3} grain sputtering cross section

Direct measurements of GCRs push towards the knee

- → Magnetic spectrometers: PAMELA, AMS-02 (, ALADINO?, AMS-100?)
- → Electromagnetic calorimeters: Fermi LAT, NUCLEON, CALET, DAMPE, ISS-CREAM?(, HERD)
- → Abundance measurements: CALET, DAMPE, SuperTIGER(, TIGERISS?)

Primary GCRs (Z > 1) show a consistent hardening at comparable rigidities

- → He (and p?) indicates subsequent softening
- \rightarrow AMS-02 demonstrates hardening in secondaries as well at lower R
- → Cause not resolved-GCR reacceleration? source classes?