Cosmic-ray Electrons at GeV—TeV energies and CALET

Nicholas Cannady ISCRA 2022 Erice, Italy

Motivation

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- Nearby sources of high-energy electrons
- Diffusion and energy loss in the Galaxy
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$$\frac{dN(E)}{dt} = D(E)\nabla^2 N(E) + \frac{\partial}{\partial E} [b(E)N(E)] + Q(E)$$

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Steady-state solution

Assuming an input power law of the form

if we find a solution for a steady state such that

we find that

or

 $Q(E) = KE^{-\alpha}$

 $\frac{dN}{dt} = 0$ and $\nabla^2 N(E) = 0$,

$$-\frac{\partial}{\partial E}[b(E)N(E)] = KE^{-\alpha}$$

 $N(E) = \frac{KE^{-(\alpha-1)}}{(\alpha-1)} \frac{1}{b(E)}$

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Ionization $\rightarrow b(E) \propto \ln(\frac{E}{m_e c^2}) \longrightarrow N(E) \approx K' E^{-(\alpha-1)}$ Bremsstrahlung $\rightarrow b(E) \propto E \longrightarrow N(E) \approx K' E^{-\alpha}$ Adiabatic cooling $\rightarrow b(E) \propto E^2 \longrightarrow N(E) \approx K' E^{-(\alpha+1)}$ Synchrotron $\rightarrow b(E) \propto E^2 \longrightarrow N(E) \approx K' E^{-(\alpha+1)}$

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Assuming injection spectrum, observed flux vs. energy tells us about significant energy loss processes $Q(E) = KE^{-\alpha}$

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$$-\frac{\partial}{\partial E}[b(E)N(E)] = KE^{-\alpha}$$

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Ionization $b(E) \propto \ln(\frac{E}{m_e c^2})$ $N(E) \approx K'E^{-(\alpha-1)}$ Bremsstrahlung
Adiabatic cooling $b(E) \propto E$ $N(E) \approx K'E^{-\alpha}$ Synchrotron
Inverse Compton $b(E) \propto E^2$ $N(E) \approx K'E^{-(\alpha+1)}$

...then for our loss term,

Beyond the accelerator, relatively little material is traversed

Synchrotron and inverse Compton dominate

$$\frac{dE}{dt} = AE^2 \qquad (A \approx 5 \times 10^{-9} Ge^{-1} yr^{-1})$$

V

Characteristic electron lifetime in this regime:

$$\frac{\Delta E}{\tau} = AE^2 \quad \xrightarrow{\Delta E \to 0.5E} \quad \tau = \frac{1}{2AE}$$

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We will not see electrons above 1 TeV that were accelerated over 1 kpc away or over 10⁵ years ago -> signatures of nearby sources

Acceleration of CR electrons

Cosmic rays up to <~ PeV energies \rightarrow diffusive shock acceleration in SNRs



X-ray synchrotron emission supports presence of high-energy electrons in SNRs



Gamma-ray observations can fit models of electronic and/or hadronic particle acceleration

Variation of parameters in models



Diffusion coefficient

Features in the spectrum above 1 TeV could confirm a nearby accelerator and constrain source and propagation model parameters

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Some dark matter candidate models include decay into leptonic channels, leading to features in the high-energy CR electron spectrum



Kaluza-Klein candidate

Fine energy resolution necessary for the detection of line signals in a continuum

Experimental results from early instruments

The observation landscape in the mid-2000's had large variations between experiments



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A feature consistently seen in the multiple flights of the ATIC detector was published in Nature in 2008 as a possible DM signature



Figure 3 | ATIC results showing agreement with previous data at lower energy and with the imaging calorimeter PPB-BETS at higher energy. The electron differential energy spectrum measured by ATIC (scaled by E^3) at the top of the atmosphere (red filled circles) is compared with previous observations from the Alpha Magnetic Spectrometer AMS (green stars)³¹, HEAT (open black triangles)³⁰, BETS (open blue circles)³², PPB-BETS (blue crosses)¹⁶ and emulsion chambers (black open diamonds)^{48,9}, with uncertainties of one standard deviation. The GALPROP code calculates a power-law spectral index of -3.2 in the low-energy region (solid curve)¹⁴. (The dashed curve is the solar modulated electron spectrum and shows that modulation is unimportant above ~20 GeV.) From several hundred to ~800 GeV, ATIC observes an 'enhancement' in the electron intensity over the GALPROP curve. Above 800 GeV, the ATIC data returns to the solid line.



Figure 4 | Assuming an annihilation signature of Kaluza-Klein dark matter, all the data can be reproduced.

Experimental results from early instruments

The observation landscape in the mid-2000's had large variations between experiments

A feature consistently seen in the multiple flights of the ATIC detector was published in Nature in 2008 as a possible DM signature

A new generation of instruments was necessary to resolve the discrepancy between experiments and to confirm or reject this very significant result



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The Calorimetric Electron Telescope (CALET)

- CALET is an ISS-borne instrument designed to study CR electrons
- Deep EM calorimeter (30 radiation lengths)
- Zenith pointing observations
- Advanced Stellar Compass star tracker for pointing knowledge
- CALET Gamma-Ray Burst Monitor (CGBM) sensitive to keV—MeV energy photons
- Deployed on Japanese
 Experiment Module Exposed
 Facility (JEM-EF)
- Operational 2015—2024+



Electromagnetic Showers

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 - Ionization
 - Bremsstrahlung
 - Annihilation (positrons)
- Photons interact with matter by
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- For sufficiently high energies, processes produce a cascade
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 - Shower ends when average particle energy too low to produce energetic secondaries



The CALET Calorimeter



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Shower imaging with the IMC – a sampling calorimeter



Showers in the TASC: a full absorption calorimeter

Fully active calorimeter: Lead Tungstate (PbWO₄)

- Dense, inorganic scintillator: lots of grammage for lots of stopping power
- Segmentation enables crude shower imaging for consistency checks with IMC reconstruction
- Fine energy deposit resolution ightarrow high fidelity electron energy reconstruction



Sample event in CALET: 250 GeV electron from flight data



CALET flight instrumentation, integration, and launch



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Calibration of the TASC on ground Asaoka et al. 2017, Astropart. Phys. 91, 1–10



Calibration of the TASC on-orbit

Temperature dependence



Proton and Helium MIP peaks compared with MC

7/31/2022

ISCRA 2022 - Cannady - CALET & Calorimetry & other time dependent corrections ³¹

Asaoka et al. 2018, Astropart. Phys. 100, 29-37

300 350 Longitude (deg)

32

Operational modes



Event trigger statistics



High-energy trigger (> 10 GeV) statistics:

- Operational time > 2,423 days^(*)
 - (*) as of May 31, 2022
- Live time fraction > 85%
- Exposure of HE trigger
 ~211 m² sr day
- HE-gamma point source exposure
 ~4.1 m² day (for Crab, Geminga)

Trigger selection

- Online trigger satisfied
- Offline trigger satisfied

Track quality cuts

- NpX > 4= && NpY >= 4
- $|NpX NpY| \le 1$
- TASC consistency

Electron selection cuts

- CHD and IMC charge selection
- Shower development threshold
- Electron likelihood
- IMC concentration
- Two-parameter K-cut



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Trigger selection



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- CHD and IMC charge selection
- Shower development thre The energy deposited in the
- **Electron likelihood**
- IMC concentration
- Two-parameter K-cut

IMC and TASC must be consistent with an electron with this total energy deposit and geometrical acceptance



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- Offline trigger satisfied

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To be discussed later... Hadronic showers have a less tightly-contained energy deposit profile



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Larger for hadronic primaries as well due to more penetrating secondaries

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- Offline trigger satisfied

Track quality cuts

- NpX > 4= && NpY >= 4
- $|NpX NpY| \le 1$
- TASC consistency

Electron selection cuts

- CHD and IMC charge selection
- Shower development threshold
- Electron likelihood
- IMC concentration
- BDT cut

Due to increasing proton contamination, a machine learning algorithm based on Boosted Decision Trees (BDT) is used to separate hadrons from electrons instead of the K-cut

BDT trained on physically-motivated parameters in different energy ranges:

- 1. The shower width in TASCX1, R_E
- 2. The fraction of energy in TASCY6, F_E
- 3. The IMCY8 shower concentration
- 4. Parameters for the TASC shower fit:
 - Shower maximum (α /b)
 - Attenuation constant (b)
 - 5% shower depth (T_{5%})
 - Goodness of fit (χ^2_{TASC})
- 5. Parameters for the IMC shower fit:
 - Exponential fit constant (p0)
 - Exponential fit slope (p1)
 - Goodness of fit (χ^2_{IMC})

Above 500 GeV, all of these except (3) are used

Selection efficiency for electrons and residual protons



FIG. 3. (a) Electron efficiency as a function of energy for each important selection step. (b) Same for protons. Black, green, yellow, blue and red histograms show the efficiency after (1) offline trigger, (2) geometrical condition with (3) track quality cuts, (4) charge selection, and (5)&(6) shower development consistency cuts, K-estimator cut with 80% efficiency, BDT cut with 80% efficiency, respectively. Adriani et al. 2017 PRL 119, 181101

Electron energy resolution

Asaoka et al. 2017, Astropart. Phys. 91, 1–10

Table 1

Summary of the error budget in the energy calibration.

MIP	Energy conversion	2.6%
Peak fitting of MC and flig Fitting range dependence Position dependence Temperature dependence Rigidity cutoff dependence Systematic uncertainty es from p/He consistency UV Laser	ght data e timated Linearity	0.6% $0.6\%^{a}$ 1.8% 1.0% $1.0\%^{a}$ 1.0% $1.4\sim 2.5\%$
Fit error APD high gain APD low gain PD high gain PD low gain		1.4% 1.5% 2.5% 2.2%
Gain Ratio	Gain range connection	1.6~ 2.1%
Fit error APD-high to APD-low gain APD-low to PD-high gain PD high to PD low gain Slope extrapolation APD-high to APD-low gain APD-low to PD-high gain PD high to PD low gain	n	0.1% 0.7% 0.1% 1.6% 2.0% ^a 1.8%
Sampling Bias		0.5% ^b

^a also considered as systematic error on energy scale

^b energy-scale systematic error only



Absolute energy scale calibration on-orbit

The Earth's magnetic field creates a low-rigidity cutoff where charged particles *cannot* reach the detector – the cutoff varies as a function of position in the orbit and direction



M. Ackermann et al./Astroparticle Physics 35 (2012) 346-353

Using a model for the magnetic field and raytracing code to determine the cutoffs, a correcting factor for the energy of electrons can be derived for individual experiments



Fig. 5. Energy spectrum of data and tracer for the McIlwain *L* interval 1.0 < L < 1.14. The black line is a fit to the data using Eq. (1) and the resulting cutoff energy has a value of 13.27 ± 0.10 GeV, and is indicated by the dot dashed line in the figure. The lower panel depicts the ratio of the two spectra.

Example figures from Fermi-LAT (AP 35 346 (2012))



Correction factor for CALET measured to be ~3.5%

Other space-borne calorimeters





Fermi Large Area Telescope (LAT)





Astropart. Phys. 95, 6 (2017)

Astrophys. J. 697, 1071 (2009)

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PHYSICAL REVIEW LETTERS 120, 261102 (2018)







Features above ~few TeV would indicate nearby source

• Electron candidate events with E >~10 TeV being studied in event-by-event analysis

Systematic errors on the CALET measurement

In addition to the statistical errors in the results, it is critical to study *systematic* errors from analysis choices

i.e. energy-dependent variation in the results can be probed via variation of analysis parameters such as (**non-exhaustive**):

- Electron survival fraction of cuts such as the K-cut, BDT, or charge selection
- MC model used for calculating instrument response
- Algorithms for reconstruction of primary particle properties such as incident direction and kinetic energy
- Binning artifacts





948.7 < E/[GeV] < 1194.3

Electron energy spectrum (preliminary update)



Modeling of contributions from CR electron sources



CR Electrons summary

- Cosmic-ray electrons are a unique channel for probing acceleration in nearby SNRs and dark matter candidates with decay via leptonic channels
- Current generation instruments are providing precision measurements of the electron spectrum into the TeV energy range
 - Calorimeters such as CALET, DAMPE, and Fermi-LAT provide direct sampling of the EM shower energy to provide accurate energy measurements
 - Magnetic spectrometers (see J. Mitchell's talk earlier today) such as AMS-02 and PAMELA offer charge-sign discrimination as well
- These current measurements are constraining structure in the spectrum, but significant questions remain
 - What systematic effects underlie the discrepancy between the flux normalization between CALET/AMS-02 and DAMPE/Fermi-LAT?
 - Are there contributions at energies above a few TeV from nearby SNRs such as Vela, Monogem, and the Cygnus loop?
 - Is the single-bin excess in the DAMPE measurement at ~1.4 a measurement at artifact or a real signal?
 - Will higher-energy measurements reveal CR electron anisotropy?