



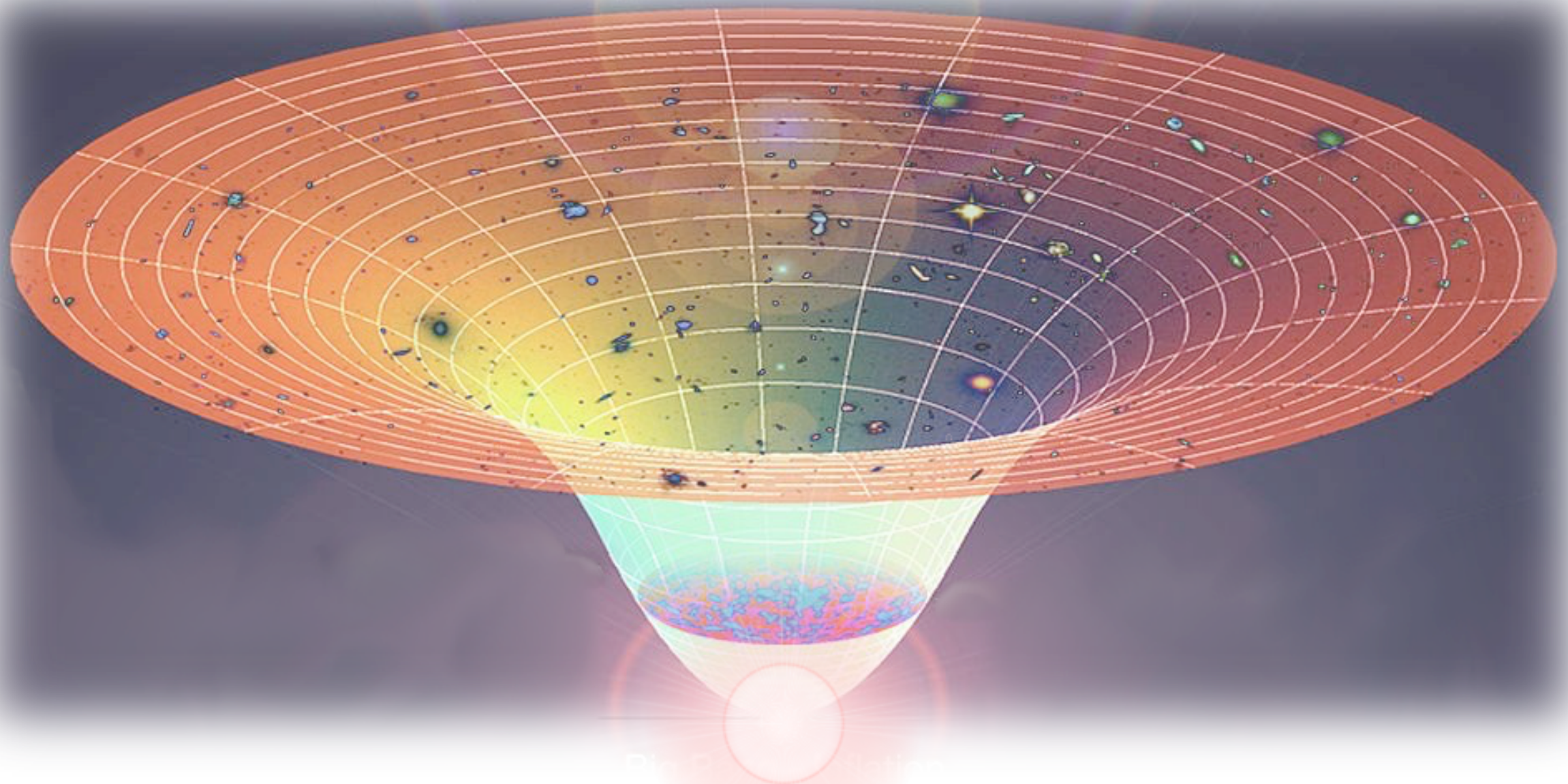
Probing gravity and dark energy in the era of multi-messenger cosmology

Alessandra Silvestri
Instituut Lorentz, Leiden U.

The standard model of Cosmology

Λ CDM: 6 parameters to describe it all!
based on General Relativity

$\{\Omega_b h^2, \Omega_c h^2, 100\theta_{\text{MC}}, \tau, \ln(10^{10} A_s), n_s\}$ measured at percent level!

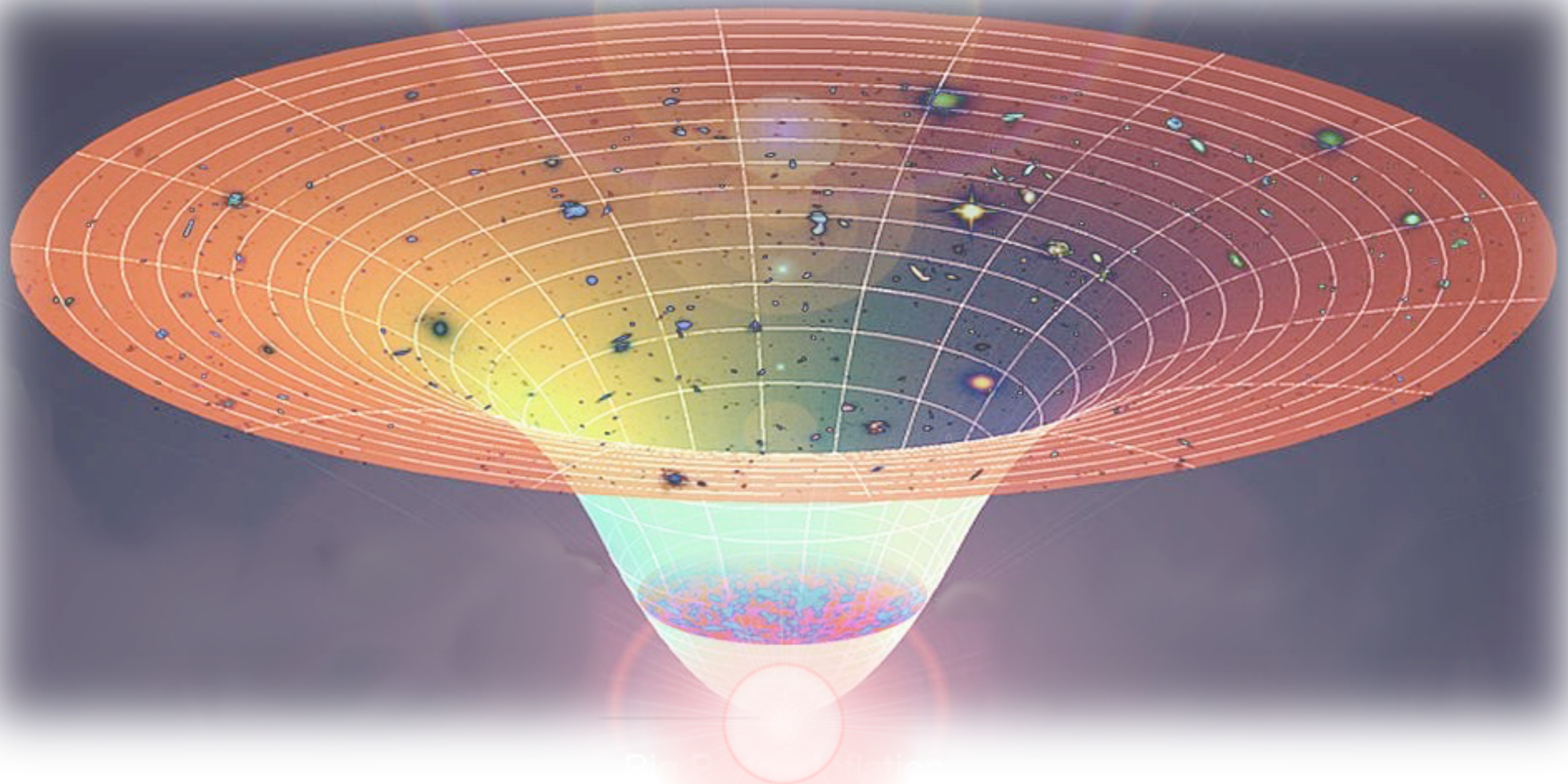


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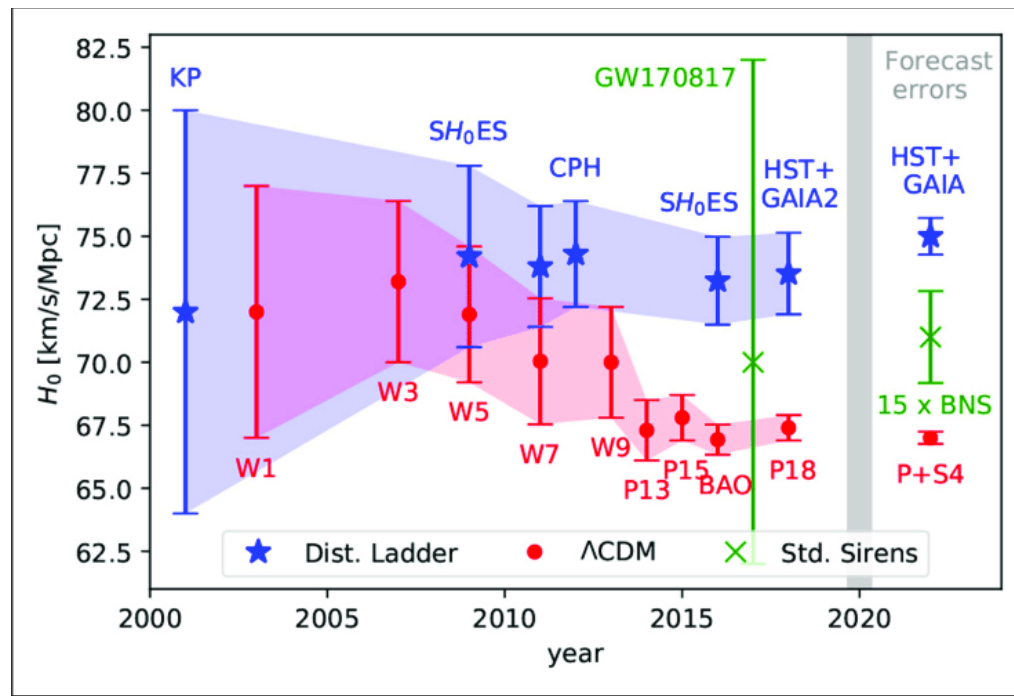
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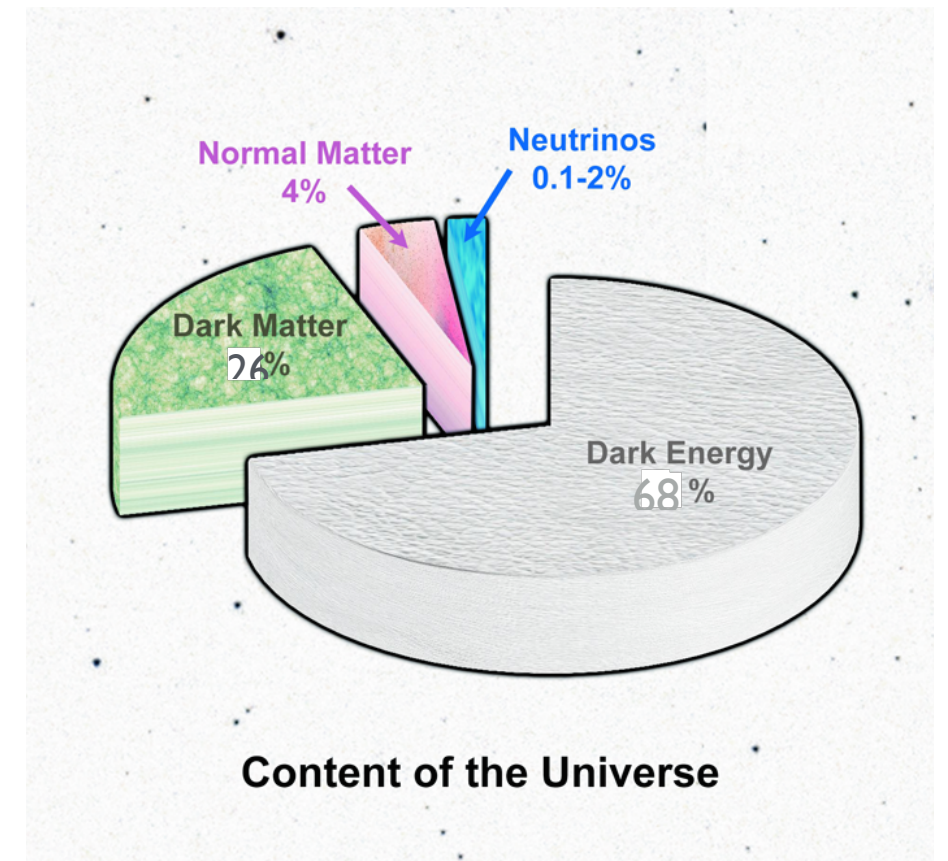


Why looking any further?

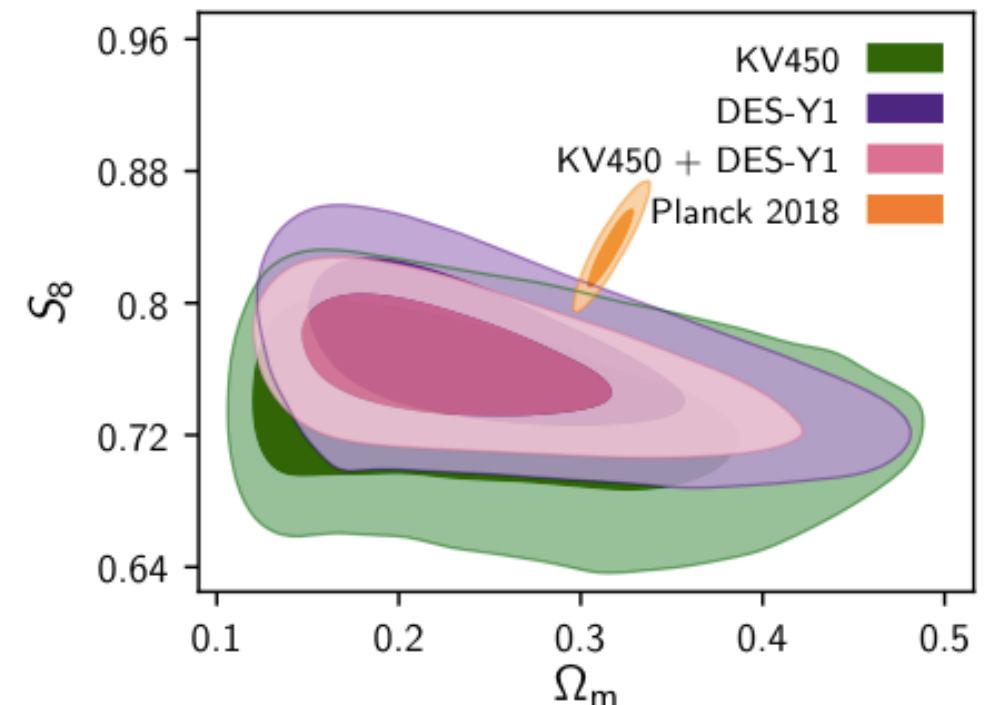


Ezquiaga et al., Front.Astron.Space Sci. (2018)

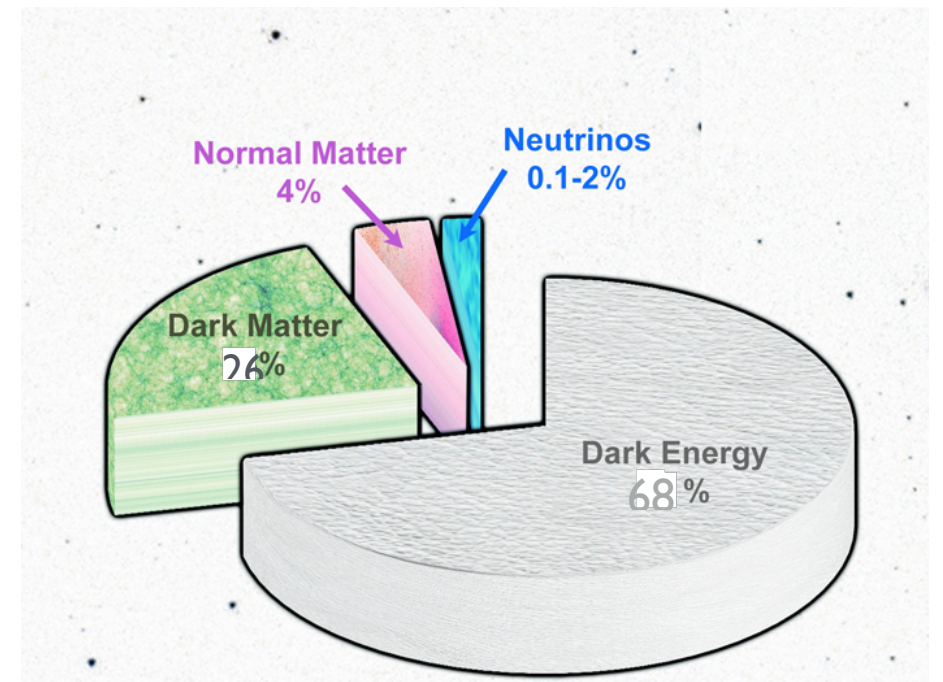
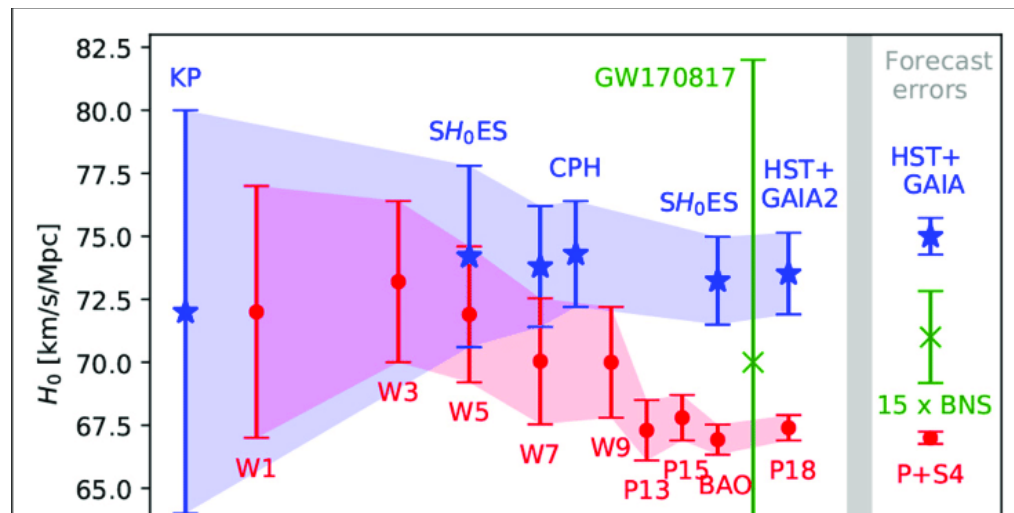
We face a mix of long standing fundamental questions (DM, DE, physics of inflation) and new tantalizing “curiosities” that make us question all aspects of our standard model as well as exploring new probes.



S.Joudaki et al., Astron.Astrophys. 638 (2020)



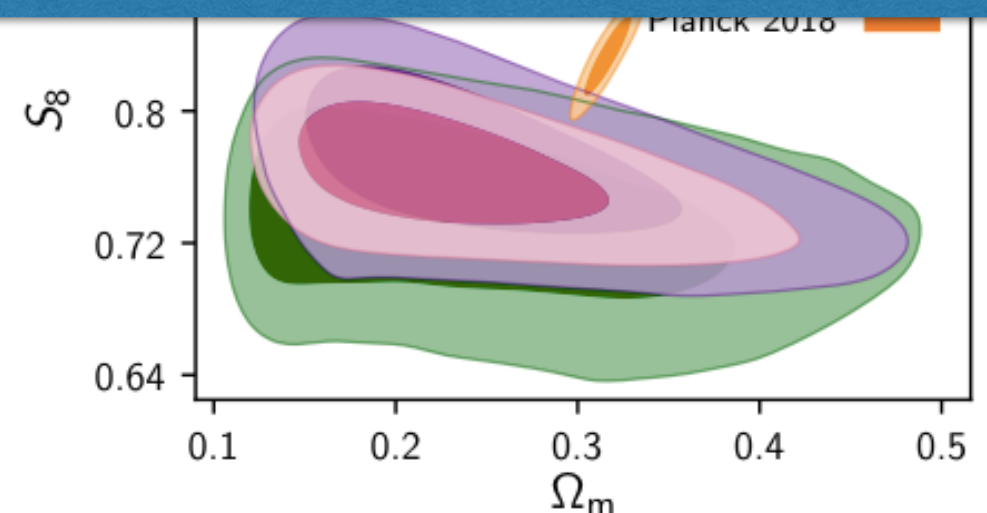
Why looking any further?



Because we can !

Upcoming surveys will provide us with a swath of data that will allow us to test gravity on cosmological scales with unprecedented precision.

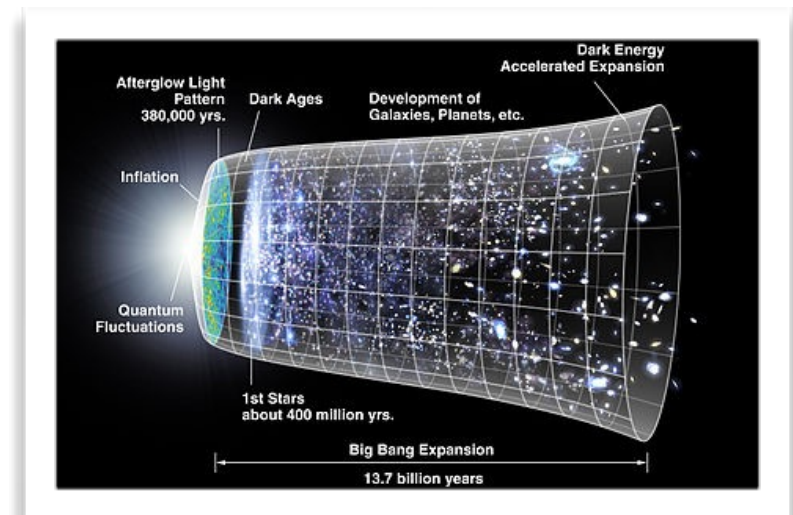
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The Standard Model of Cosmology

Einstein eqs. applied to the background FLRW metric of our Universe $ds^2 = -dt^2 + a^2(t)d\vec{x}^2$ give the Friedmann eq.:

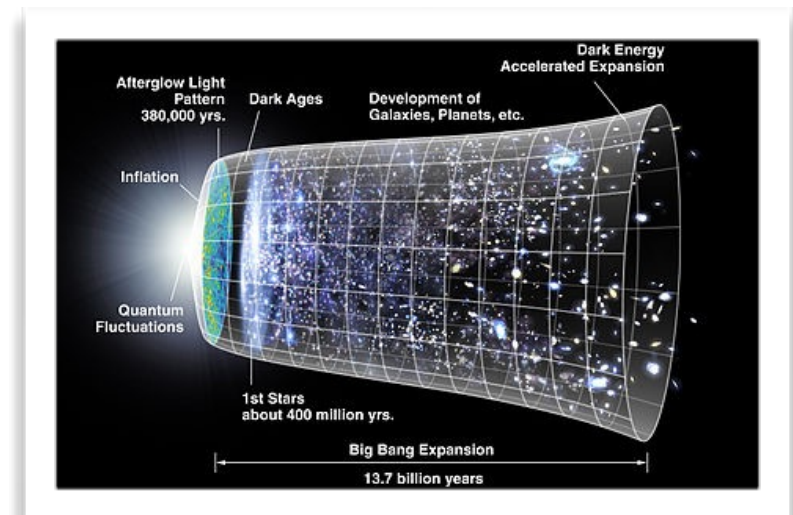
$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{1}{3M_P^2} \left(\frac{\rho_m^0}{a^3} + \frac{\rho_r}{a^4} + \rho_\Lambda^0 \right)$$



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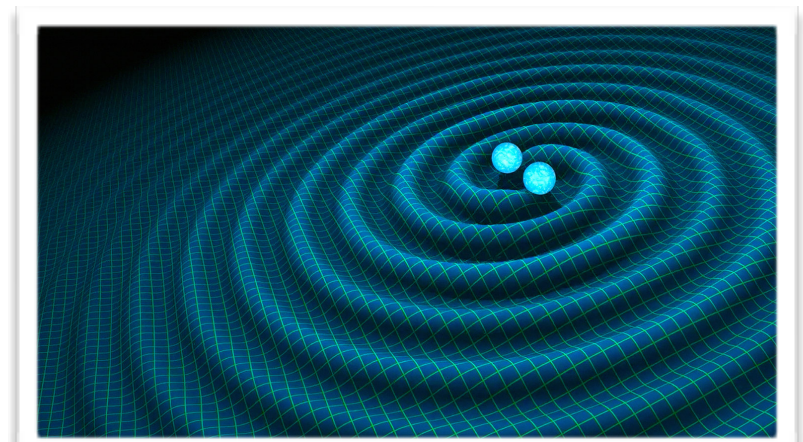
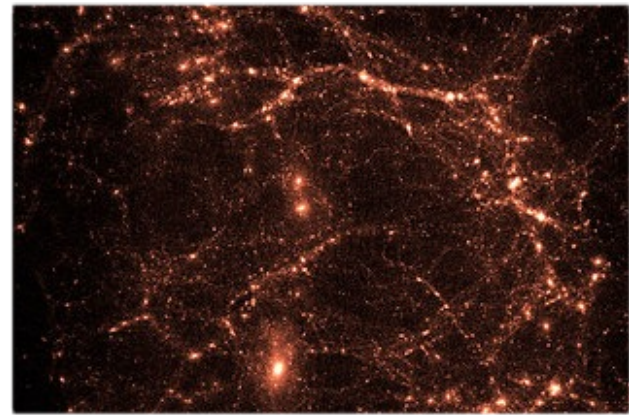
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If we now include **perturbations**

$$ds^2 = -[1 + 2\Psi(t, \vec{x})]dt^2 + a^2(t)[1 + 2\Phi(t, \vec{x})]d\vec{x}^2 + h_{ij}(t, \vec{x})dx^i dx^j$$

the Einstein equations at **linear** level are **decoupled**:



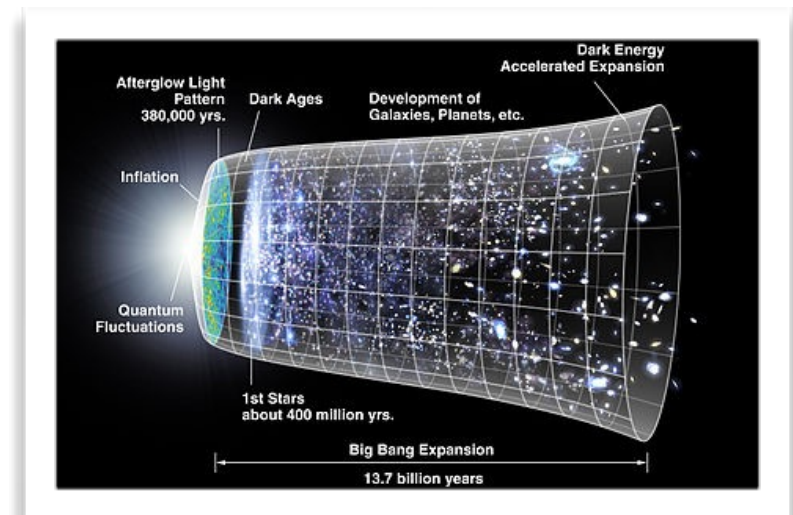
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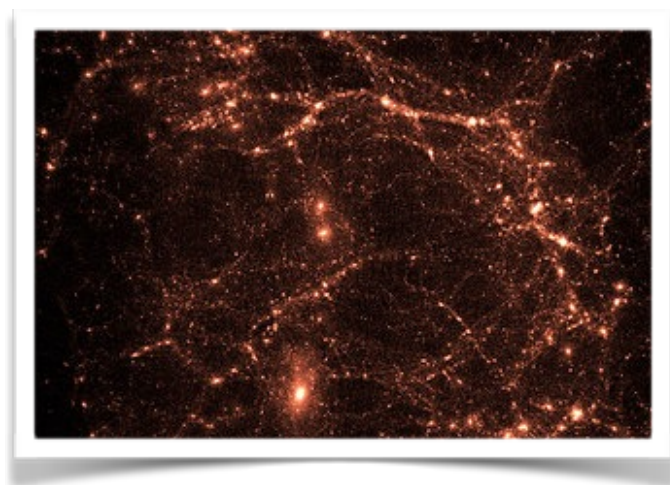
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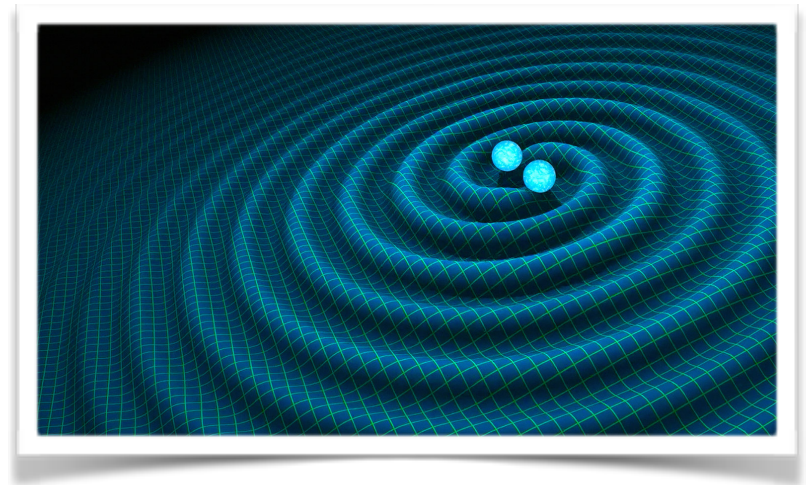
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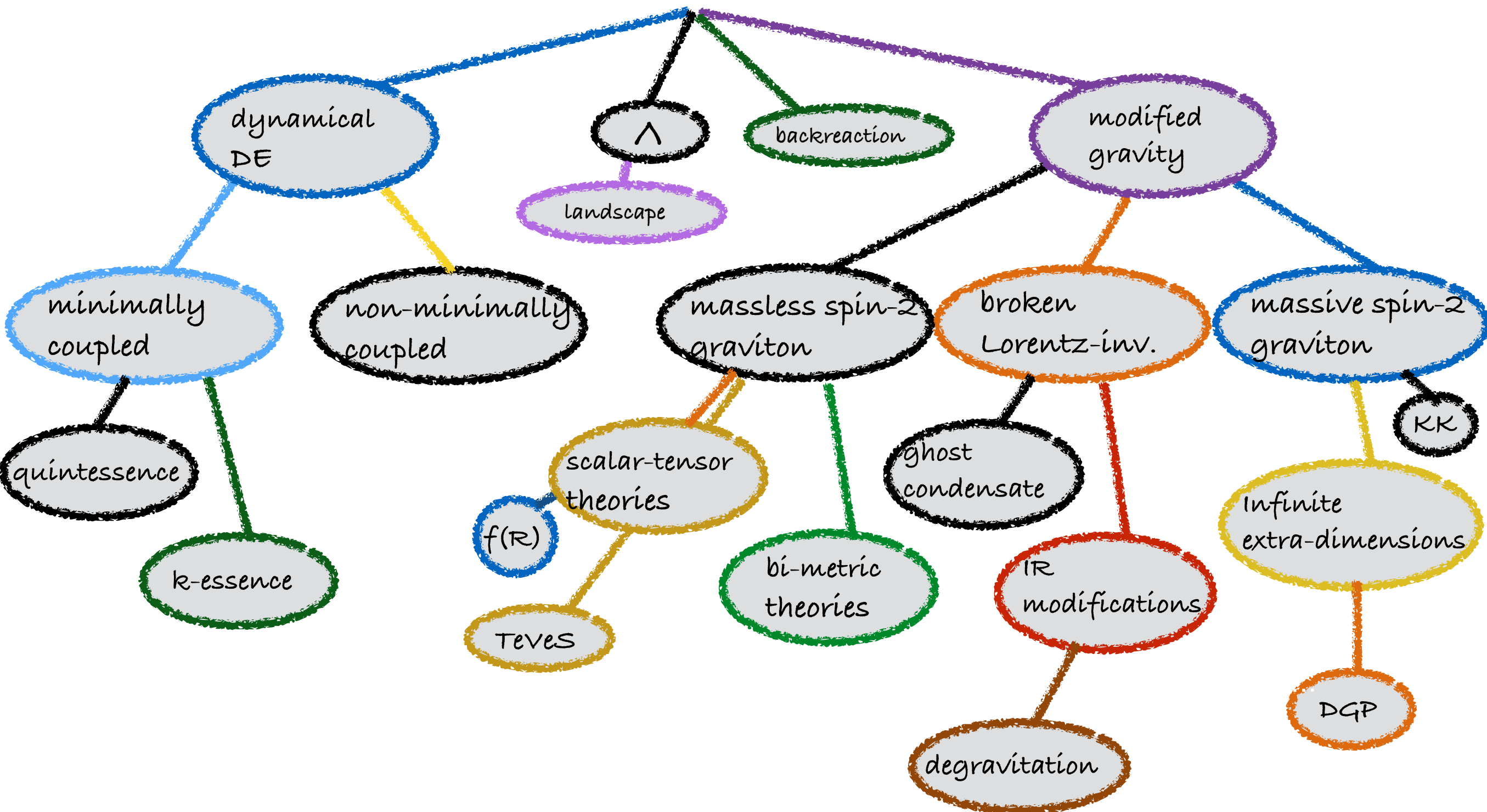
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Beyond LCDM ?

Let me focus on dark energy and modifications of gravity

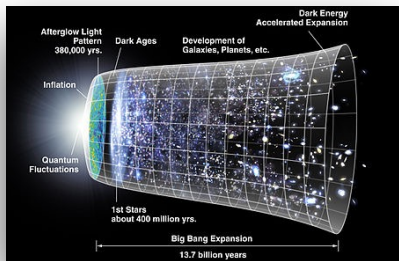
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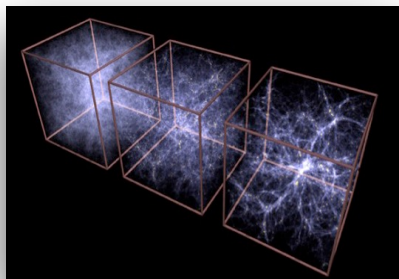
Beyond LCDM - LSS frontier

The equations for (linear scalar) perturbations become significantly more complicated. Yet we can capture the effects in few phenomenological functions:



Expansion:

$$\frac{H^2}{H_0^2} = \frac{\Omega_r}{a^4} + \frac{\Omega_M}{a^3} + \Omega_{DE} a^{-3} \int (1 + w_{DE}(a))$$



Clustering:

$$k^2 \Psi = -\mu(a, k) \frac{a^2}{2M_P^2} \rho \delta$$



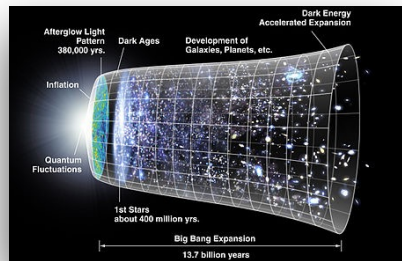
Lensing:

$$k^2 (\Phi + \Psi) = -\Sigma(a, k) \frac{a^2}{M_P^2} \rho \delta$$

for scalars only!

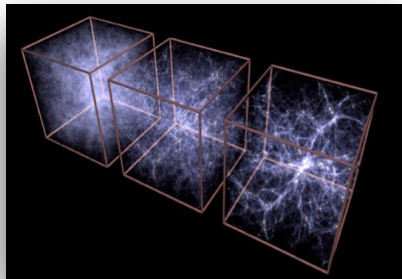
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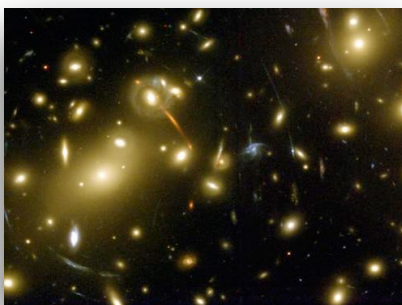
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Ongoing and upcoming wide field imaging and spectroscopic redshift surveys are in line to map more than a 100 cubic-billion-light-year of the Universe: exquisite measurements of **expansion rate** & reconstruction of cosmic structure **growth rate** and **lensing** to $\sim O(1)\%$ over wide redshift range.

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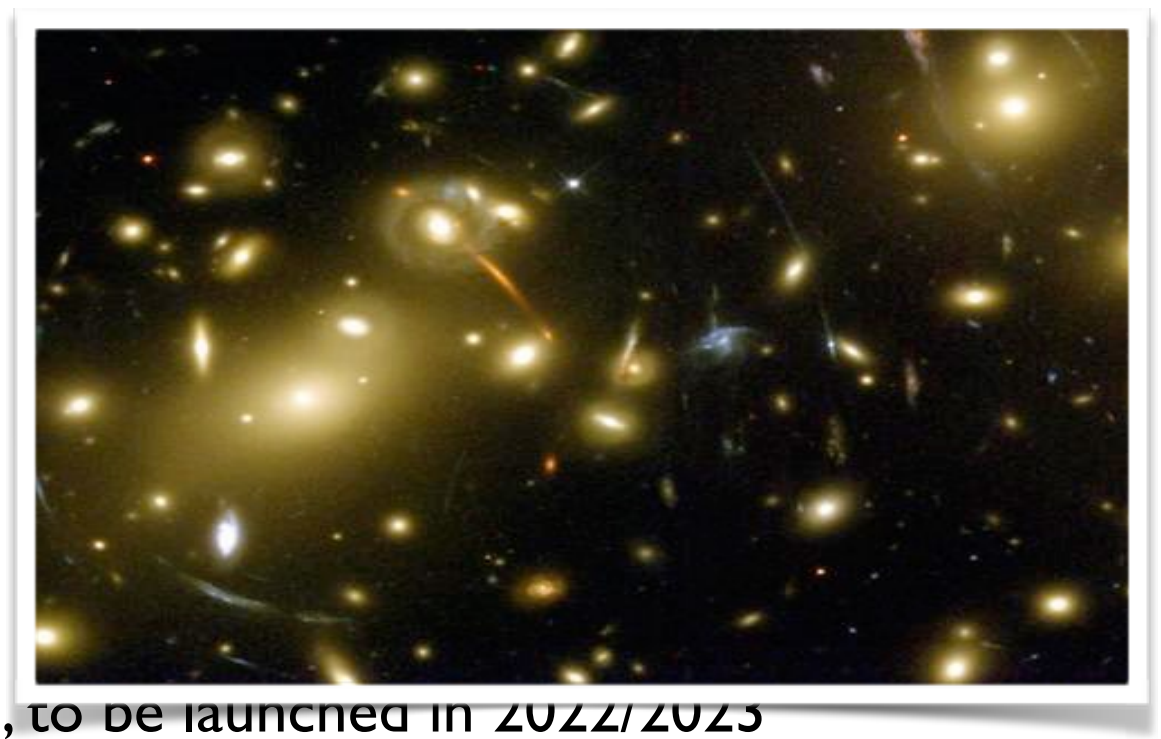
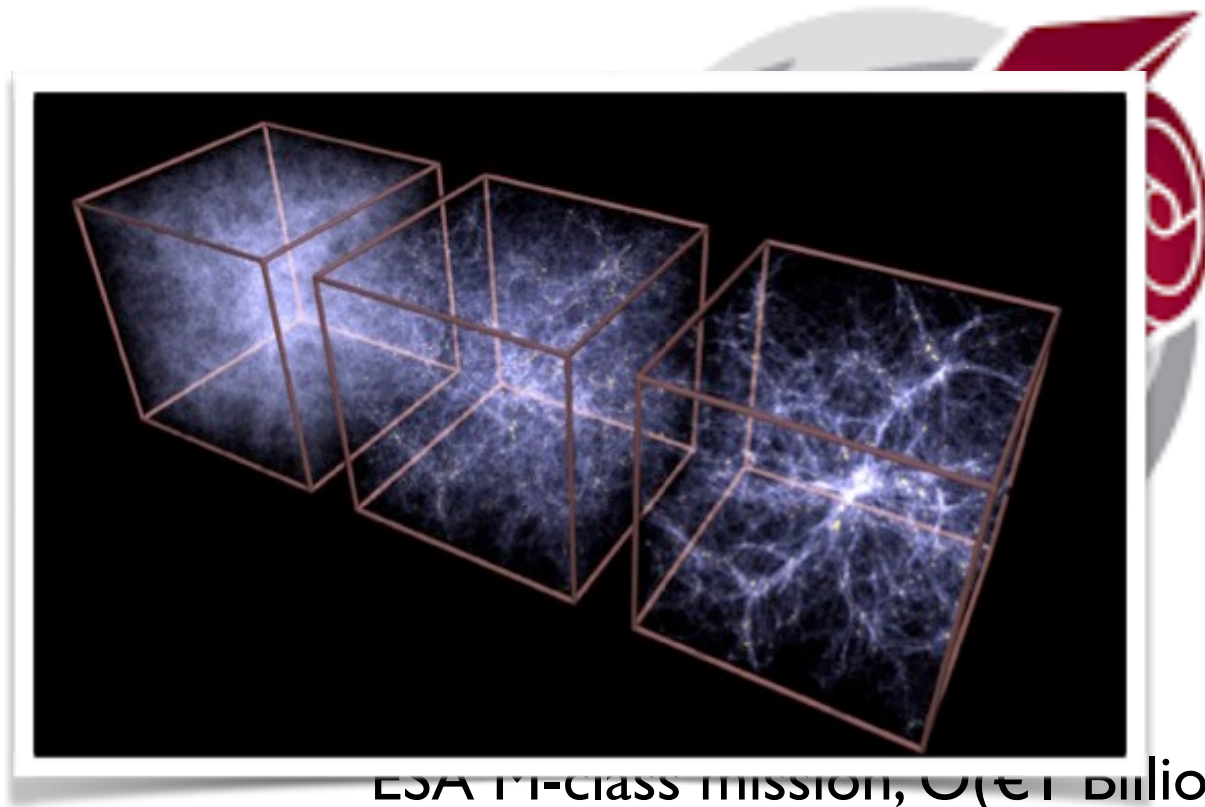


ESA M-class mission, O(€1 Billion), to be launched in 2022/2023

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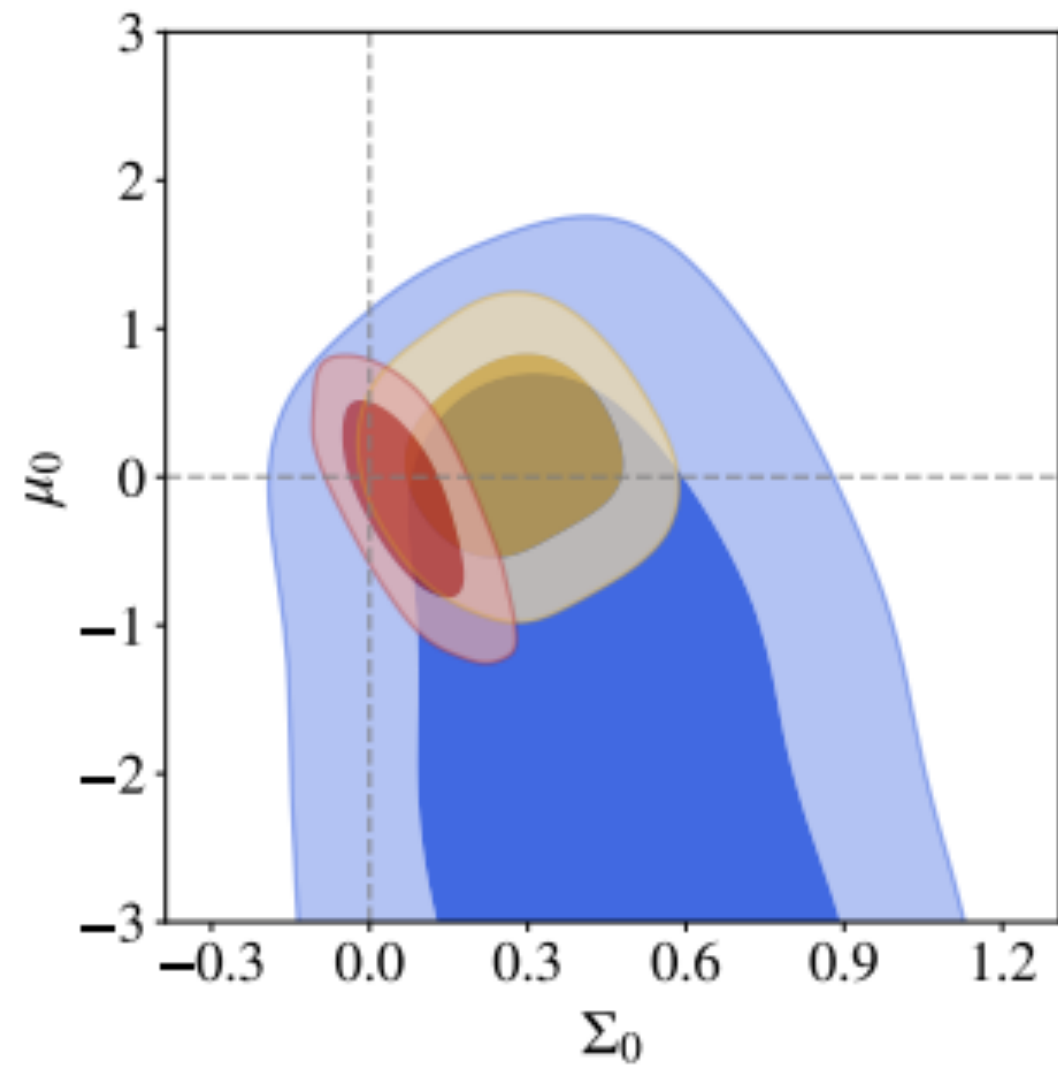
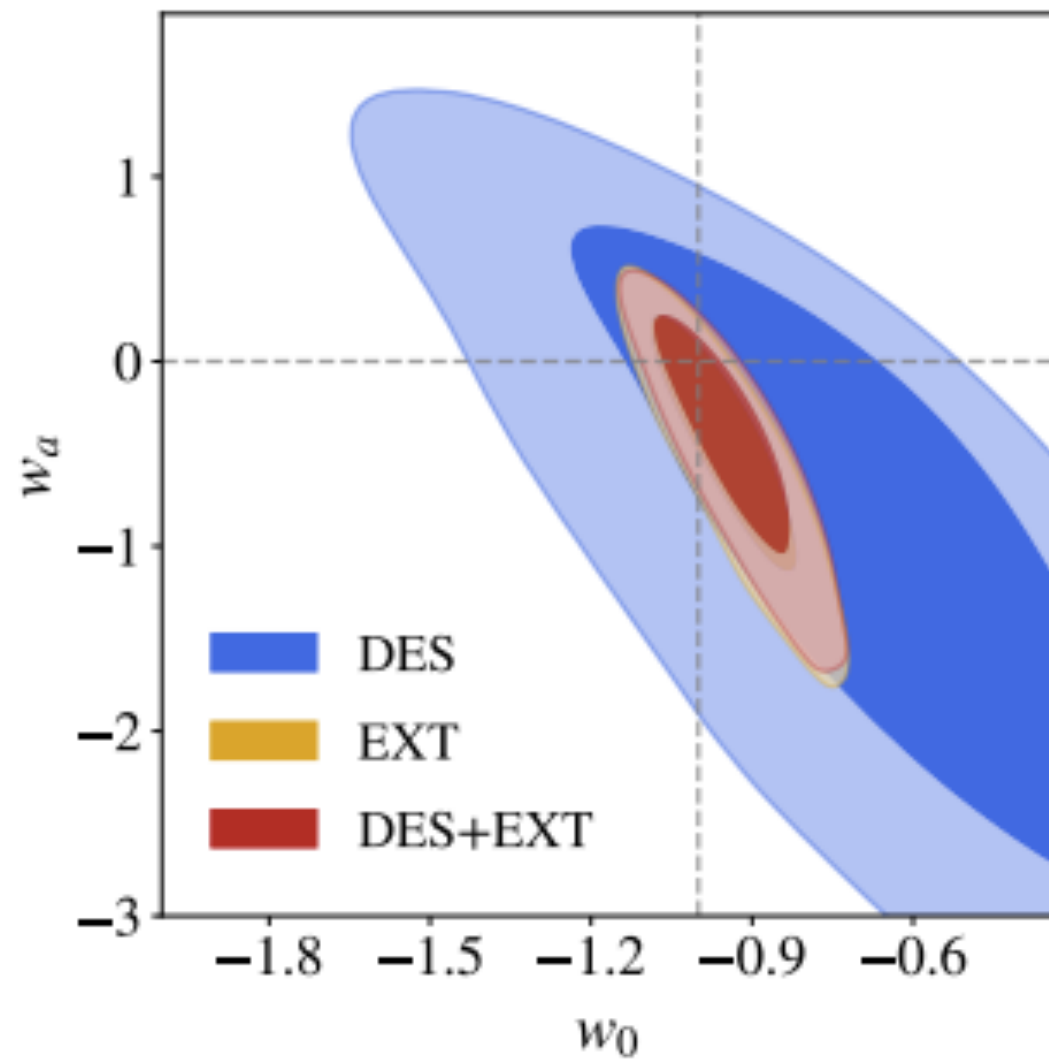


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$$w_{\text{DE}} = w_0 + w_a(1 - a)$$

$$\Sigma(a) = \Sigma_0 \frac{\Omega_\Lambda(a)}{\Omega_\Lambda}$$

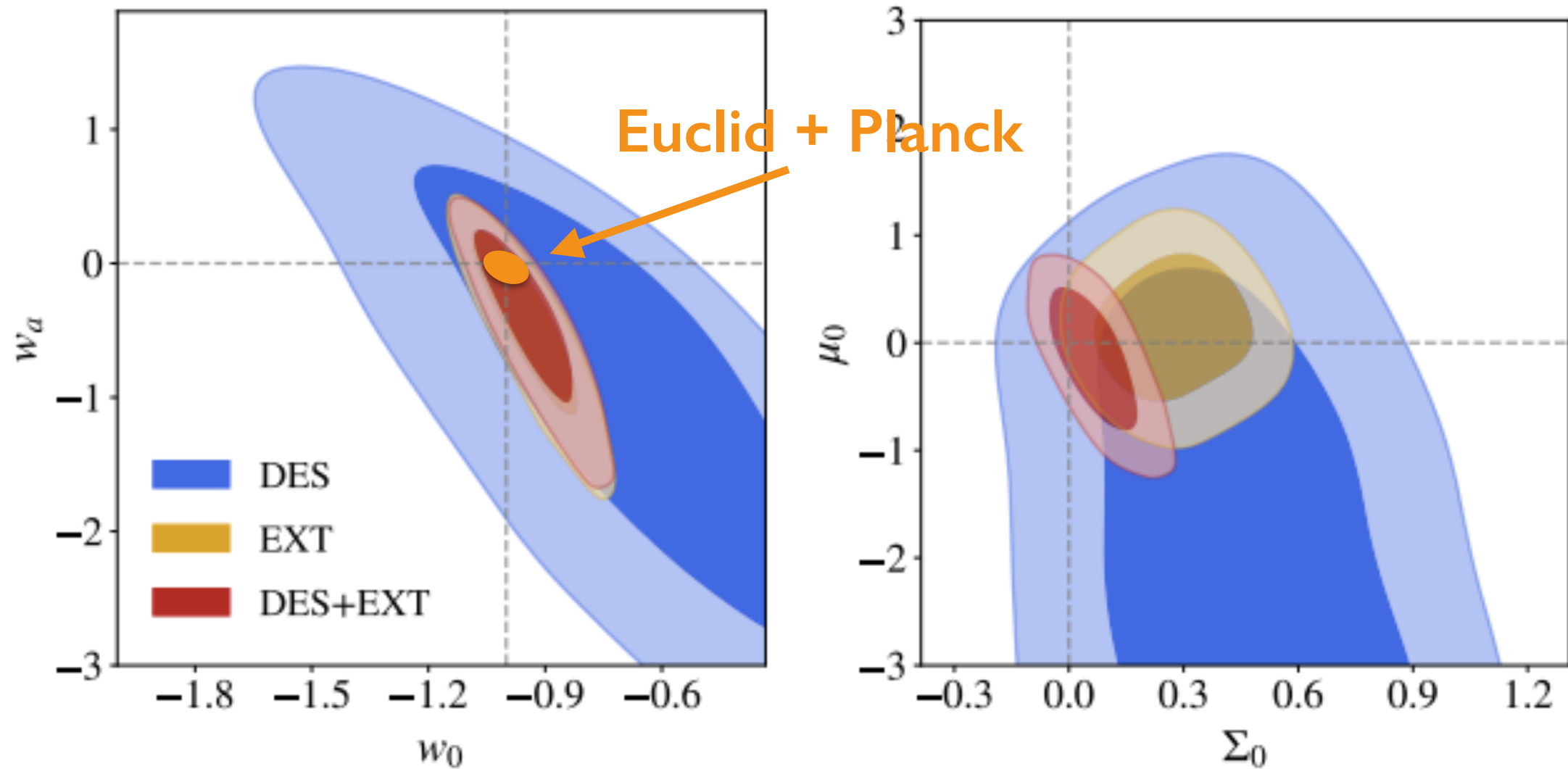


DES year 1

Beyond LCDM - LSS frontier

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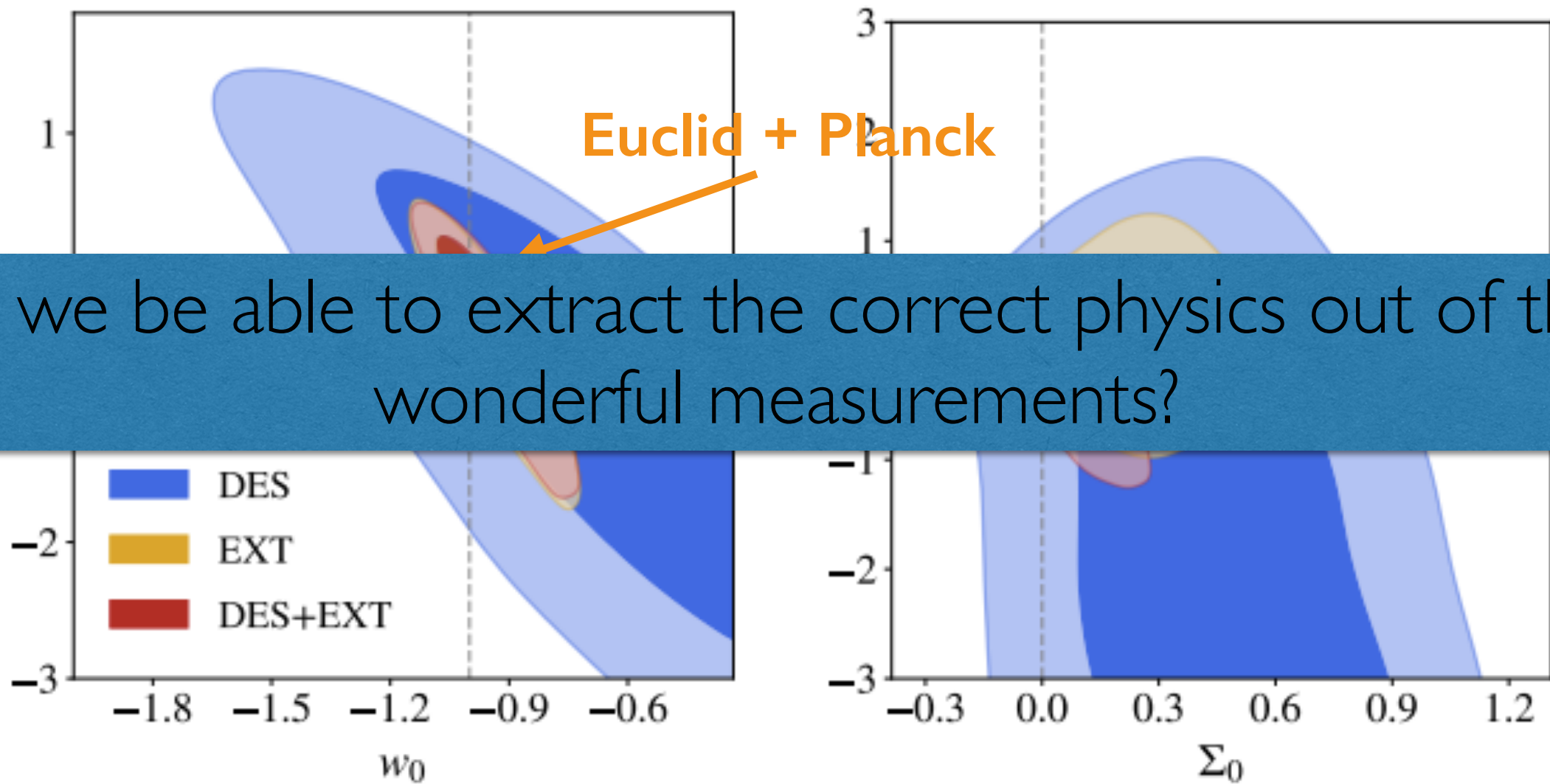


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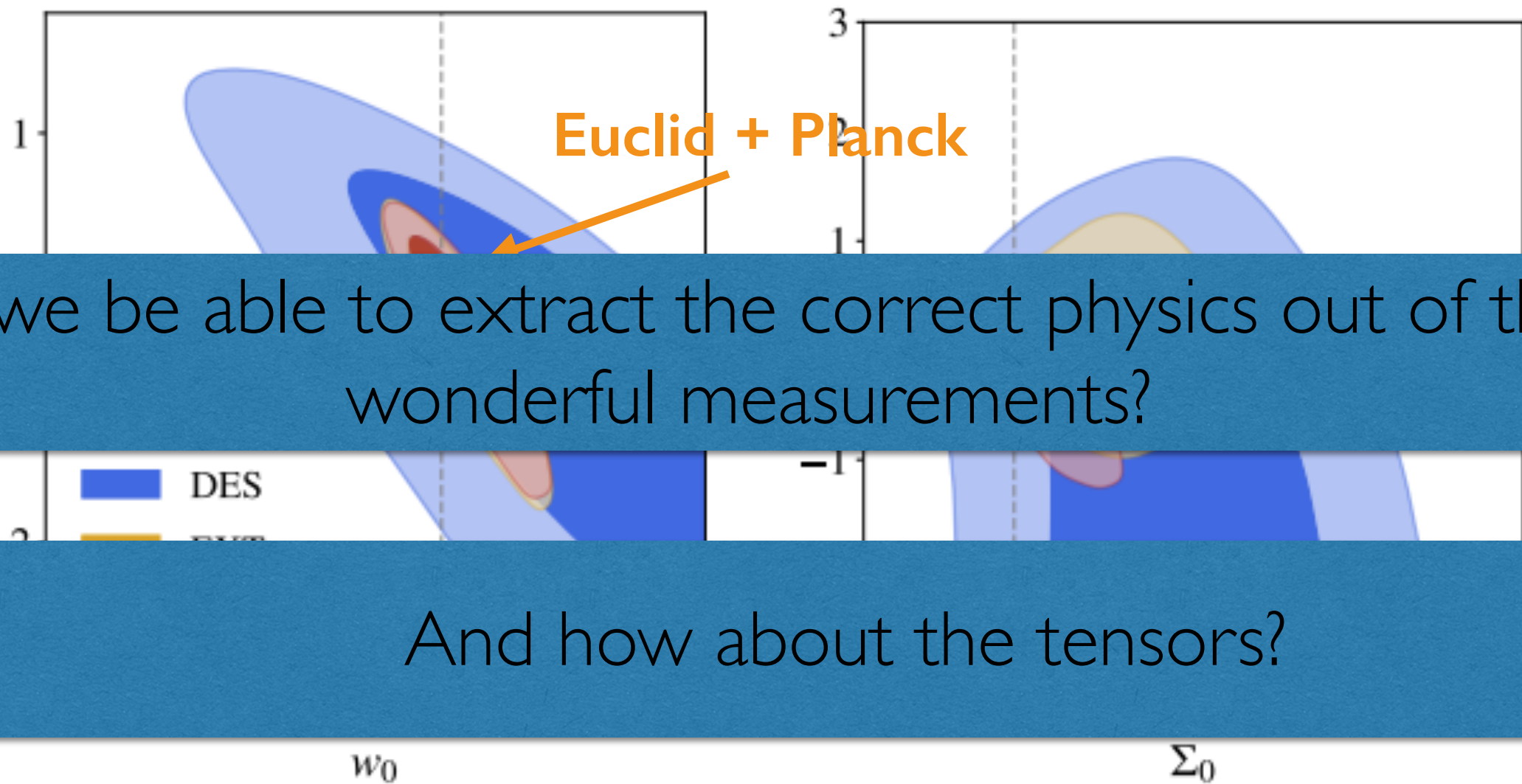


Will we be able to extract the correct physics out of these wonderful measurements?

Beyond LCDM - LSS frontier

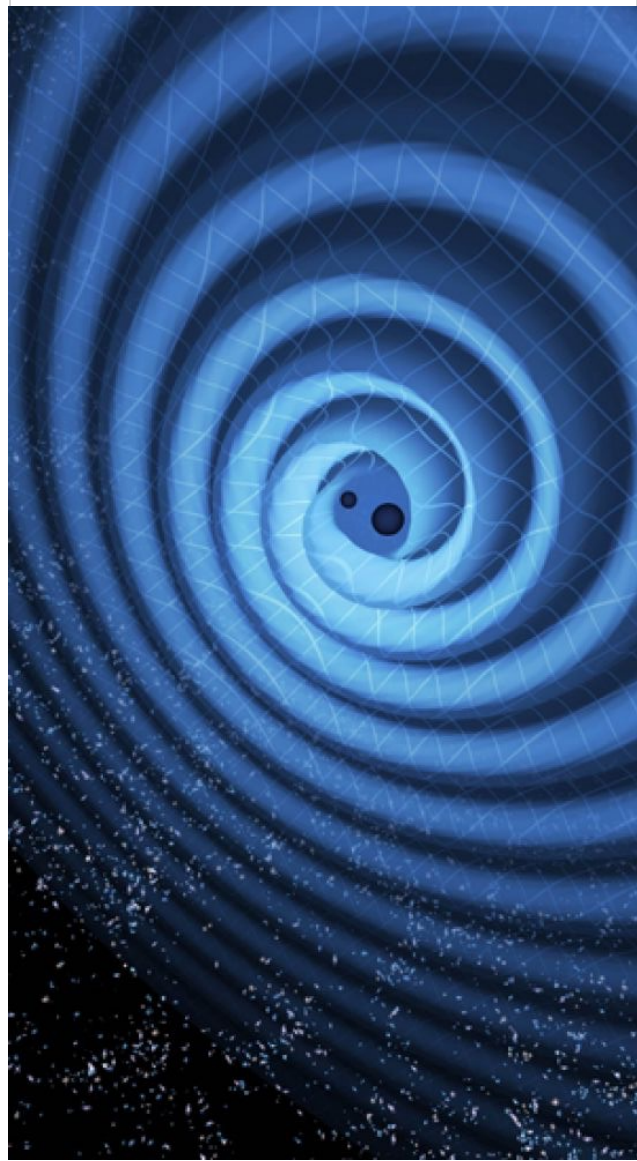
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GW170817

& GRB170817A



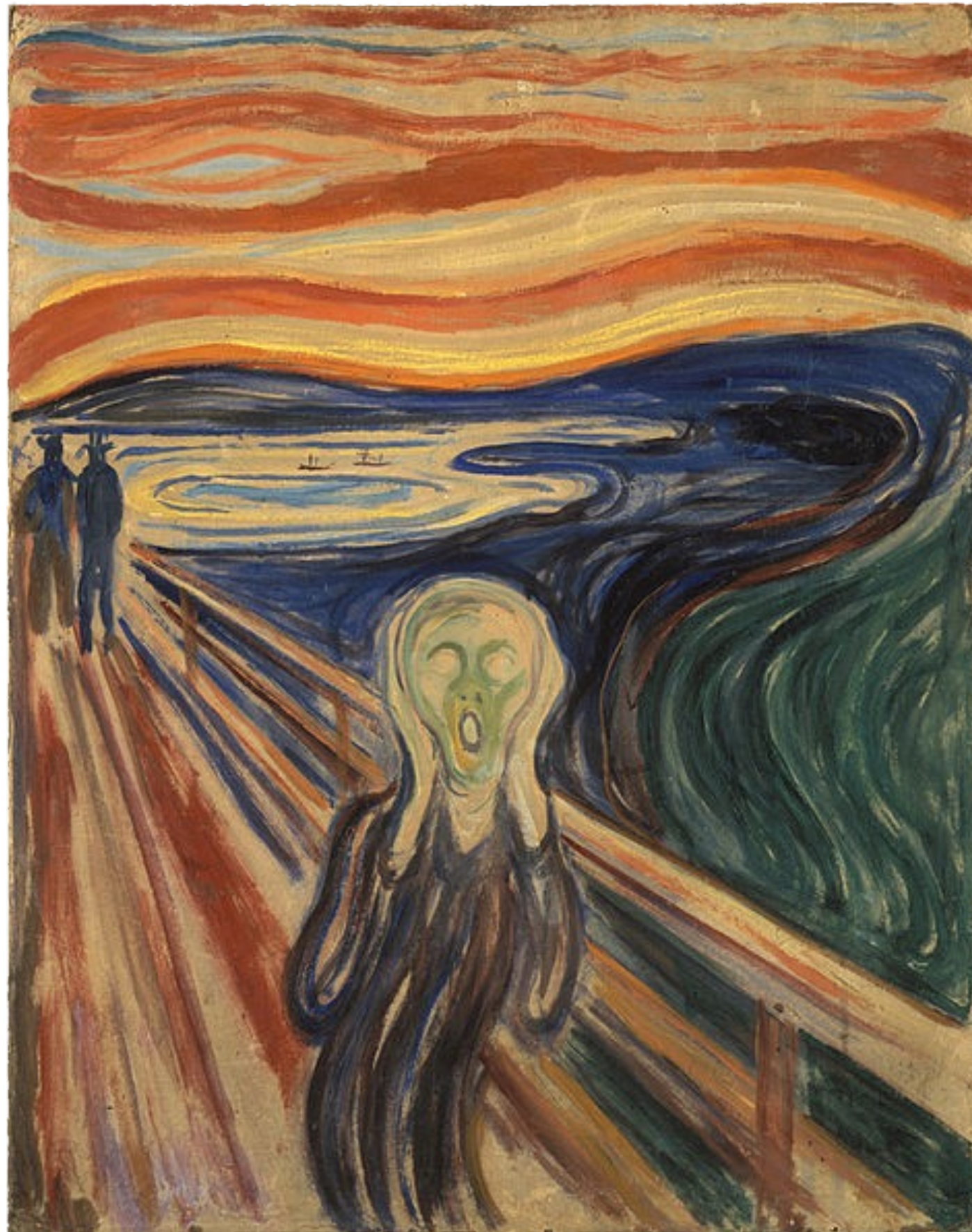
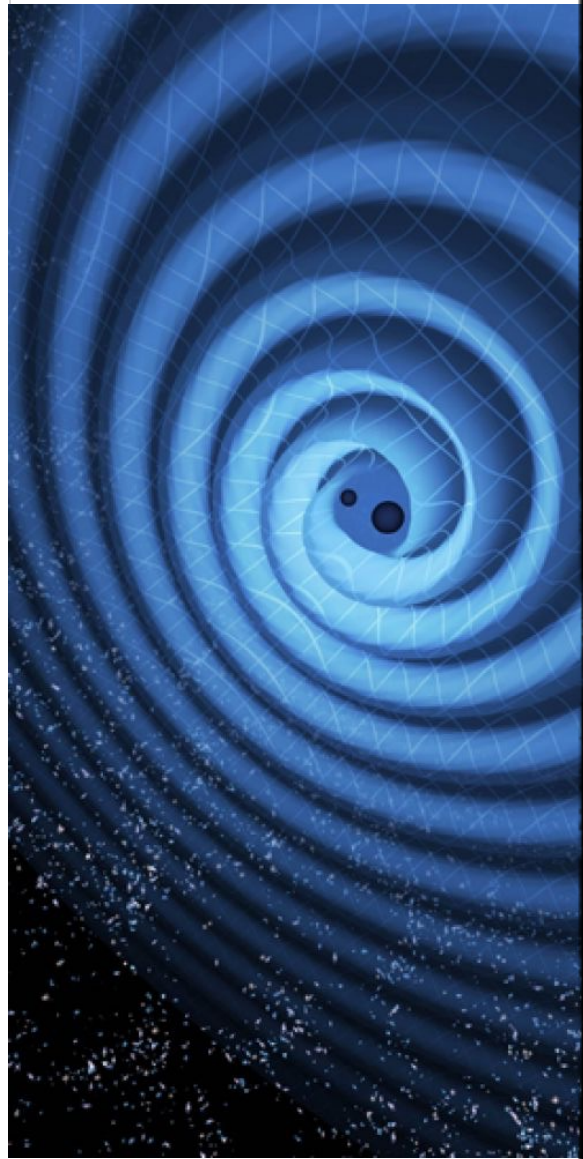
Using the observed **time-delay** btw GRB and GW,
(1.74 ± 0.05)s , they placed very stringent limits on the speed of gravity:

$$-3 \cdot 10^{-15} < \frac{c_T^2}{c^2} - 1 < 7 \cdot 10^{-16}$$

and just like that, many modified gravity models were ruled out* !

GW170817

170817A



The scream, E. Munch

GW,

h the speed of gravity:

were ruled out* !

Will we be able to extract the correct physics out of these wonderful measurements?


And how about the tensors?

Effective Field Theory of Dark Energy

A unified action for linear perturbations

$$S^{(2)} = \int dx^3 dt a^3 \frac{M_*^2}{2} \left\{ \delta K_{ij} \delta K^{ij} - \delta K^2 + (1 + \alpha_T) \delta_2 \left(\sqrt{h} R / a^3 \right) + \alpha_K H^2 \delta N^2 + 4\alpha_B H \delta K \delta N \right\}$$

$\alpha_M \equiv H^{-1} \frac{d \ln M_*^2}{dt}$



- α_M Running of Planck's constant, generated by non-minimal coupling
- α_T Deviation of speed of GWs from unity; non-zero whenever there is a non-linear derivative coupling of the scalar field to the metric. Same non-linearity is responsible for non-zero anisotropic stress.
- α_K Quantifies the independent dynamics of the scalar-field
- α_B Signals a coupling between the metric and the scalar-field

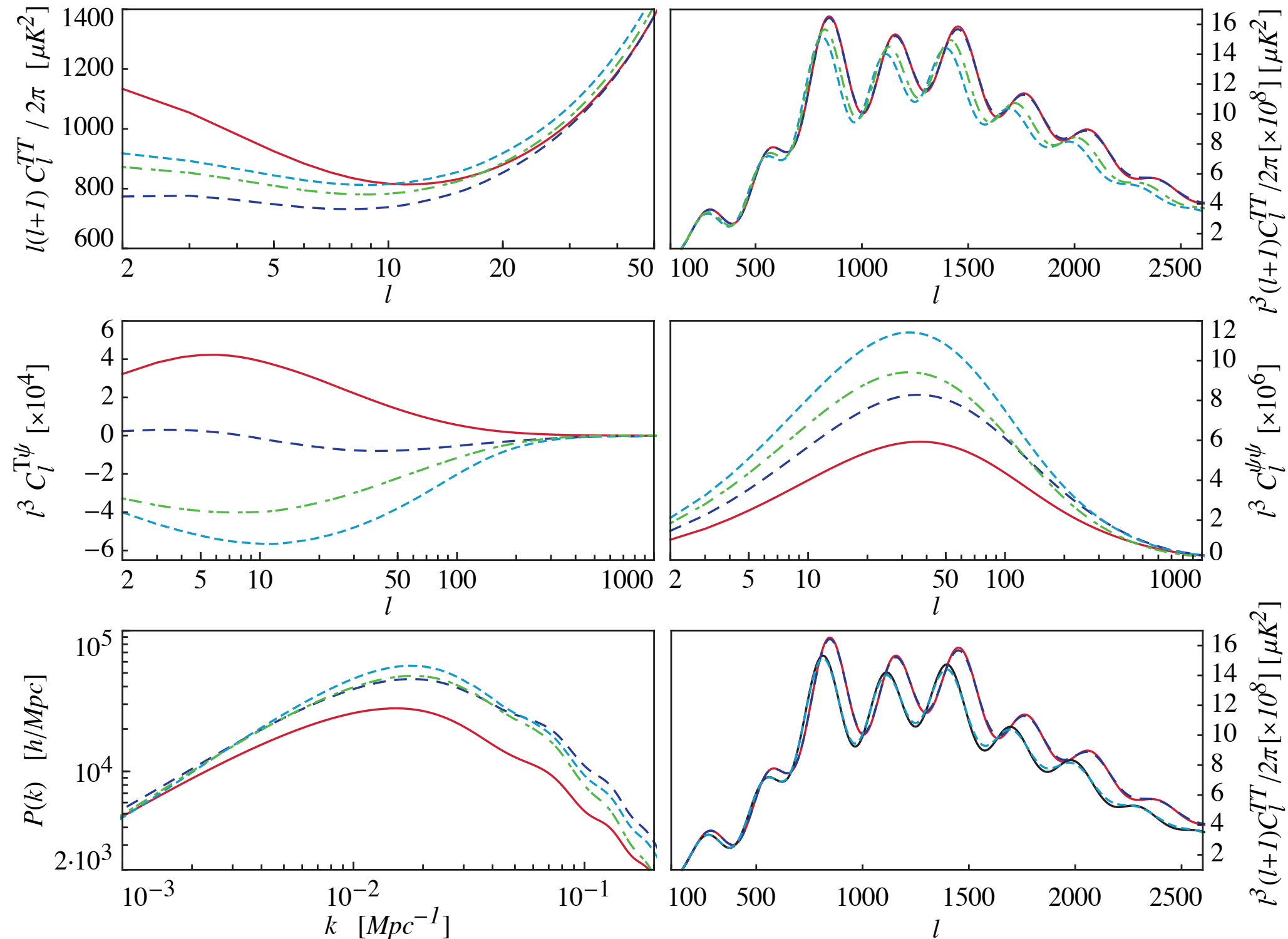
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$\alpha_M \equiv H^{-1} \frac{d \ln M_*^2}{dt}$

- Guided by symmetry principles
- Unified and physically informed framework
- Unified treatment of vast range of observables

LSS in Beyond LCDM with EFTCAMB



Gravitational Waves Phenomenology

From the same unified action, we can now study also the propagation of tensors:

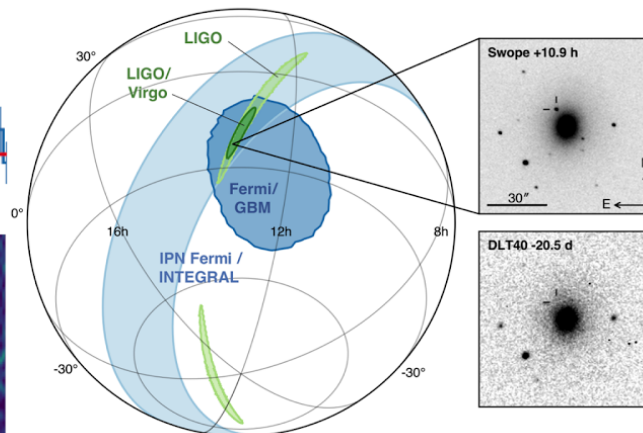
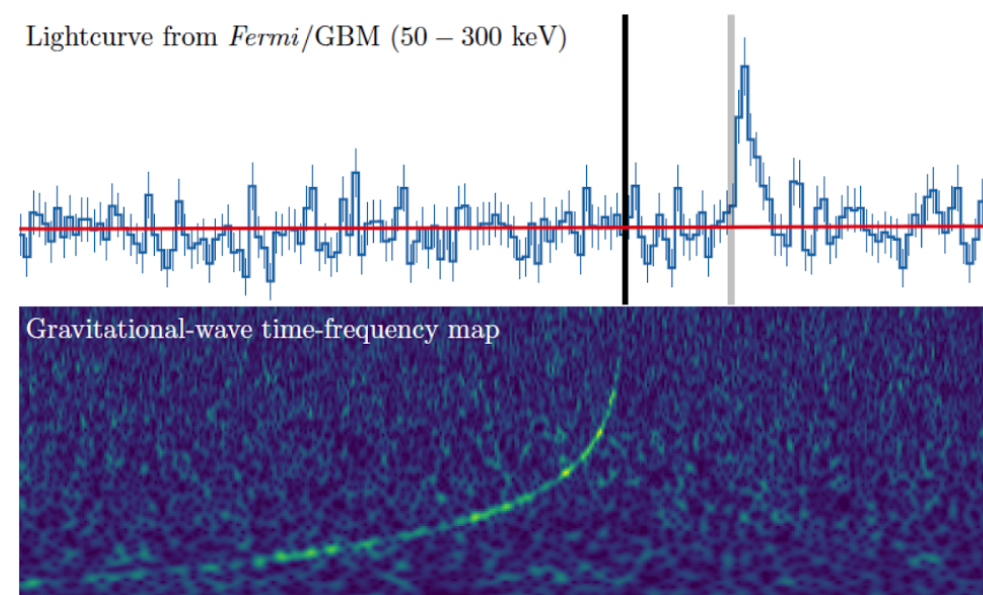
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GW170817 + GRB170817A



$$-3 \cdot 10^{-15} < \frac{c_T^2}{c^2} - 1 < 7 \cdot 10^{-16}$$

$$\alpha_T = 0$$

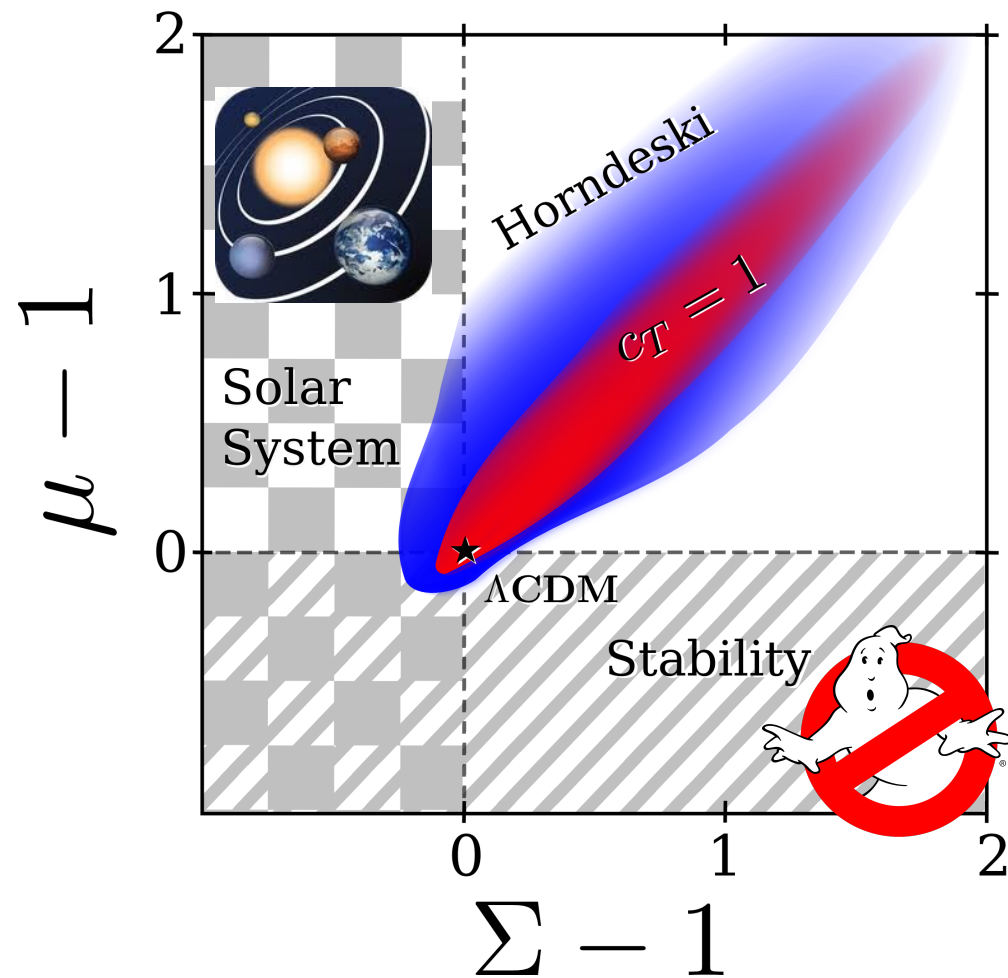
This cuts a big chunk of interesting, self-accelerating models

Creminelli & Vernizzi, PRL 2017
 Ezquiaga, Zumalacarregui, PRL 2017
 Baker et al., PRL 2017

Modified Gravity after GW170817

& GRB170817A

- A modified propagation of GW and a gravitational slip $\mu \neq \Sigma$ are intertwined



More specifically, a standard speed of propagation implies:

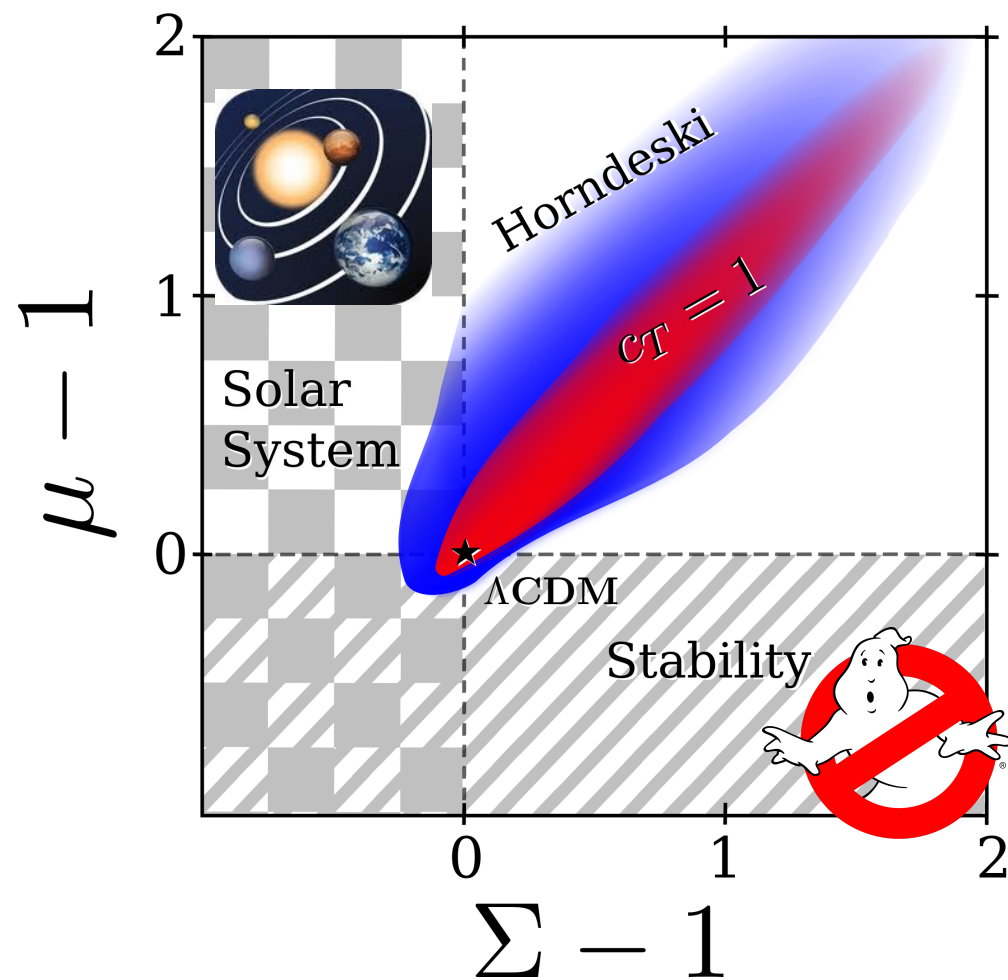
$$\Phi(k \ll Ma) = \Psi(k \ll Ma)$$

No gravitational slip on large scales for scalar-tensor theories!

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Euclid will measure this slip with ~ 1 -10% accuracy.

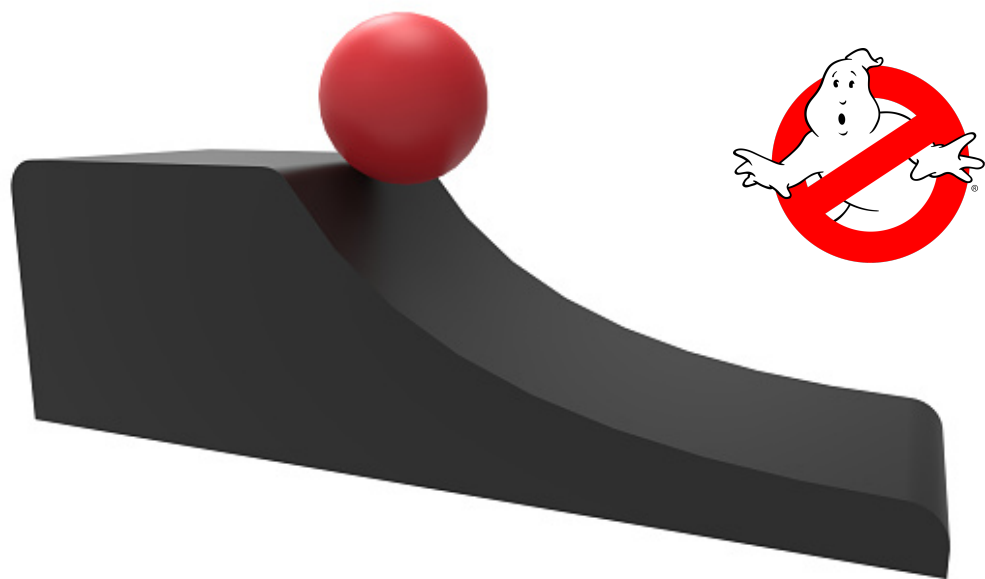
Were we to find a non-null value, then we would detect some “beyond scalar-tensor”.

Theoretical Priors

Expanding the given action up to second order in the perturbations, and removing spurious DOFs, we can inspect the dynamics of perturbations:

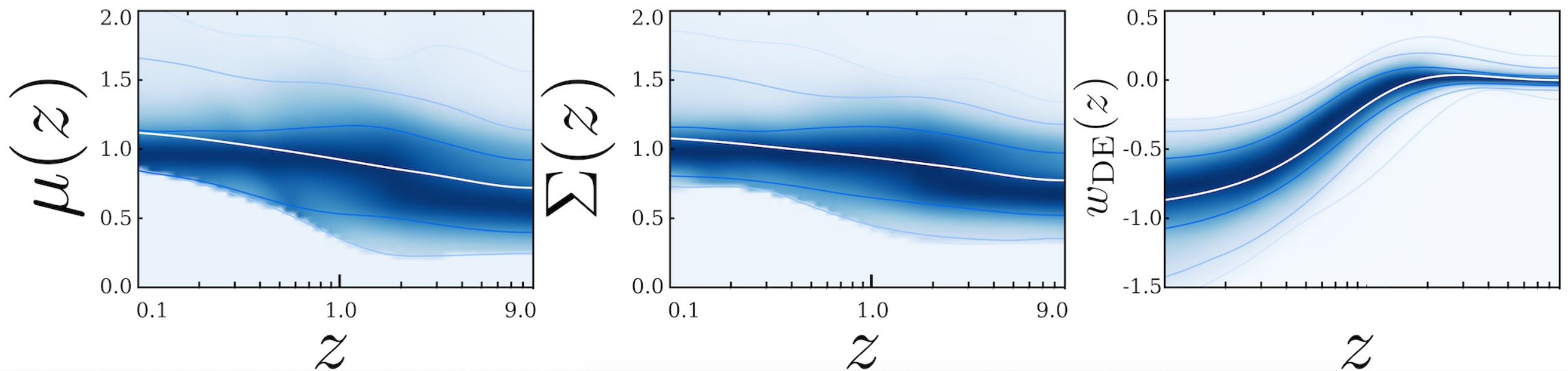
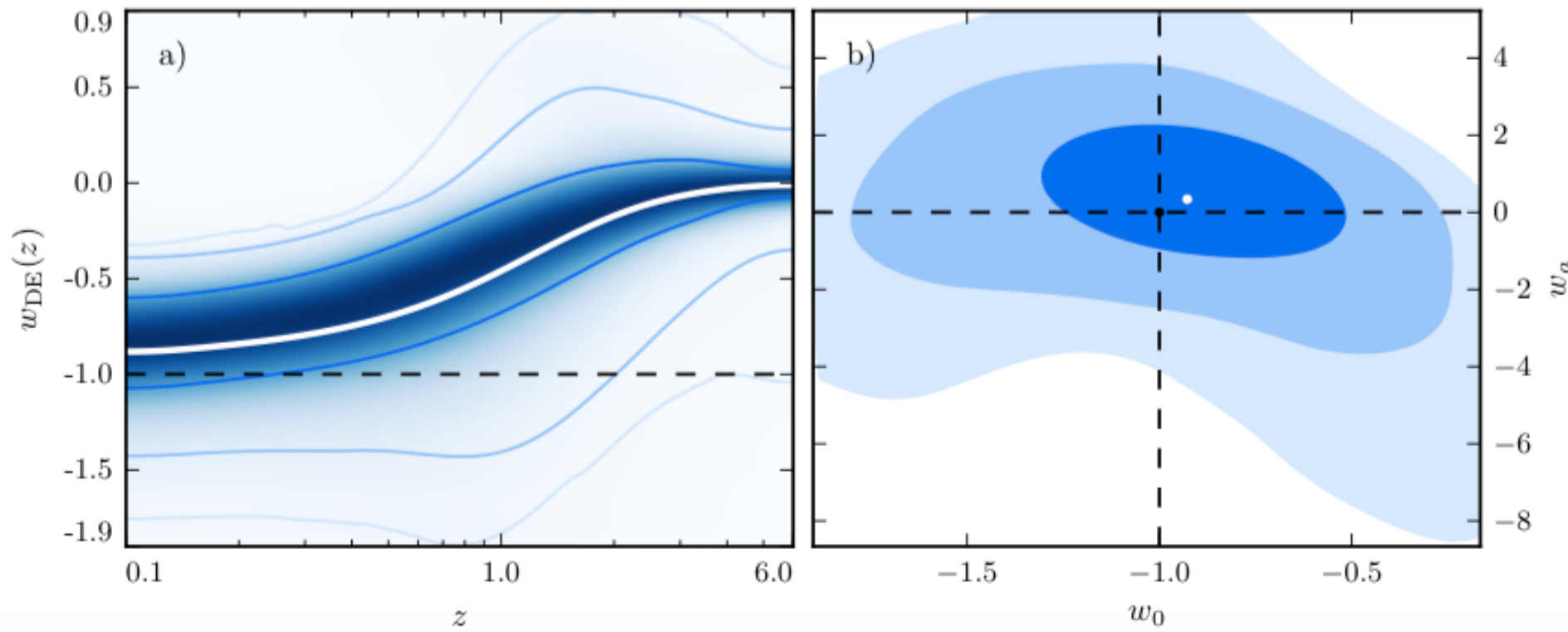
$$S_{\xi, h^T}^{(2)} = \frac{1}{(2\pi)^3} \int d^3k dt a^3 \left\{ \left[\mathcal{L}_{\xi\xi} \dot{\xi}^2 - k^2 G \xi^2 \right] + \frac{A_T}{8} \left[\left(\dot{h}_{ij}^T \right)^2 - \frac{c_T^2}{a^2} \left(h_{ij}^T \right)^2 \right] \right\}$$

and the avoidance of instabilities translate into a set of precise conditions on the free functions:

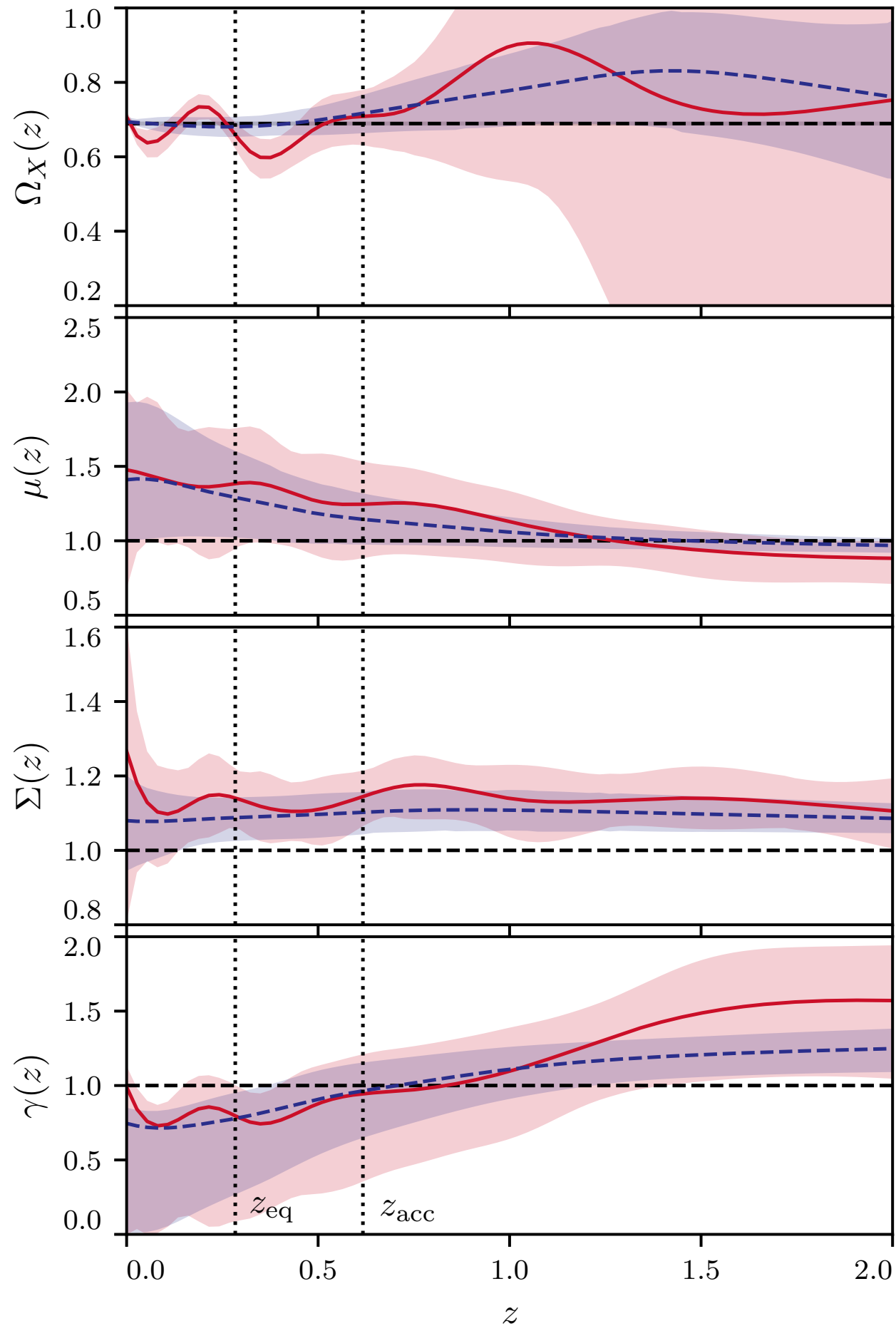


$$\begin{aligned} \mathcal{L}_{\xi\xi} &> 0 \\ c_s^2 &\equiv \frac{G}{\mathcal{L}_{\xi\xi}} > 0 \\ A_T &> 0 \\ c_T^2 &> 0 \end{aligned}$$

Theoretical Priors for LSS



Reconstruction with theoretical priors



Pogosian, Peirone, Zhao, Li, Raveri, Koyama, AS,
arXiv:2107.12990
arXiv:2107.12992

First joint non-parametric reconstruction
of gravity from Large Scale Structure data,
under guidance of theoretical viability
conditions.

#weak lensing

DES

#Planck 2015 TT, lowTEB, lowl

#BAO Likelihoods

BAO: 6DF, MGS, eBOSS DR12 (in 9 bins by YTW), eBOSS DR14 (a
z_eff by GBZ),

eBOSS DR12 Lyman alpha

#Others:

Pantheon

WiggleZ_MPK

GW & beyond LCDM

So, we are left with:

$$\ddot{h}_{ij} + (3 + \alpha_M) H \dot{h}_{ij} + (1 + \cancel{\alpha_T}) k^2 h_{ij} = 0$$

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then, inside the horizon

$$h \propto \frac{1}{\tilde{a}} \quad \frac{\dot{\tilde{a}}}{\tilde{a}} = \frac{3 + \alpha_M}{2} H$$

$$d_L^{\text{GW}}(z) = \frac{a(z)}{\tilde{a}(z)} d_L^{\text{EM}}(z)$$

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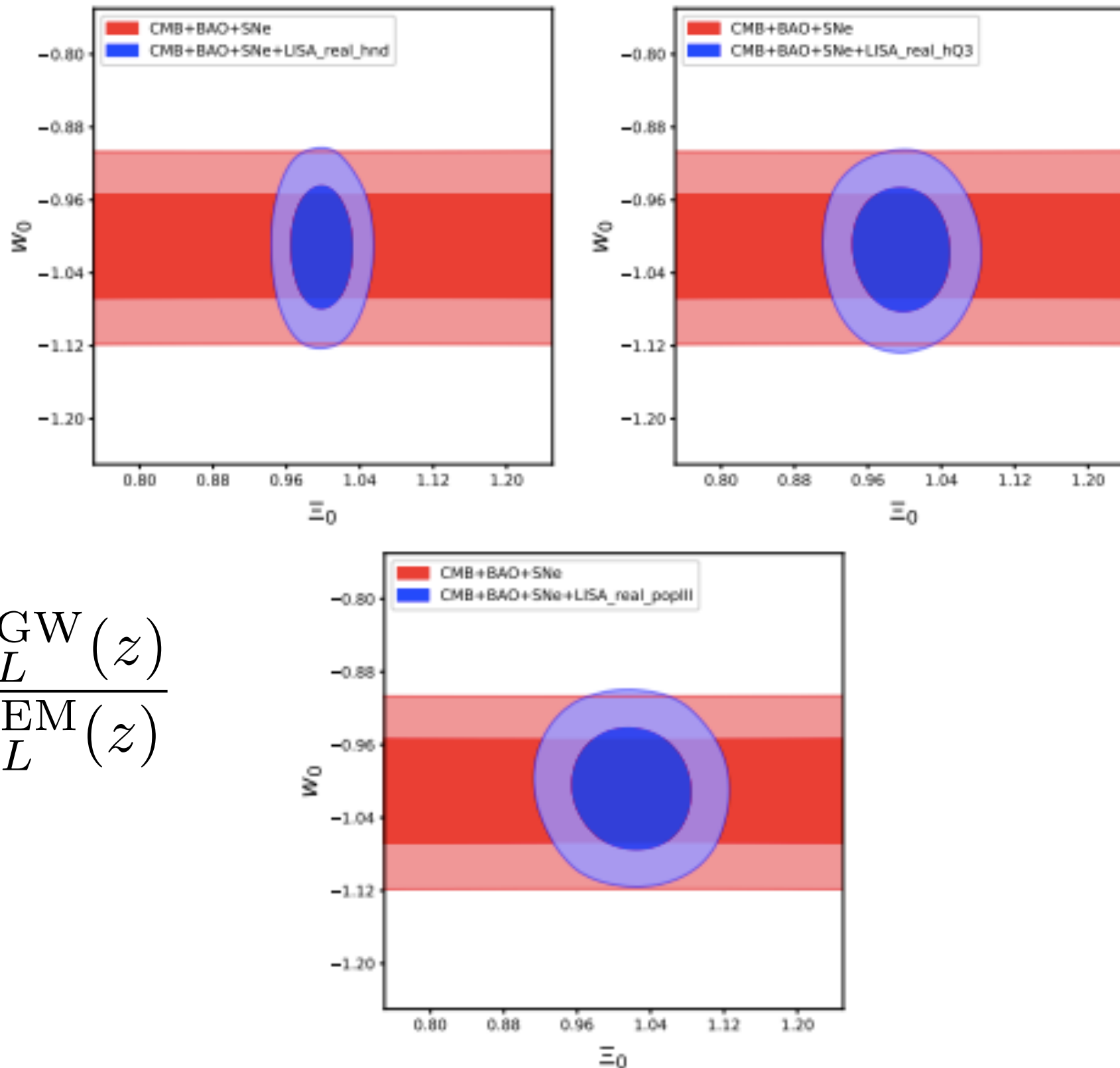
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$$h \propto \frac{1}{\tilde{a}} \quad \frac{\dot{\tilde{a}}}{\tilde{a}} = \frac{3 + \alpha_M}{2} H \quad (1 + z) \int_0^z \frac{cdz'}{H(z')}$$

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GW & beyond LCDM



$$\Gamma(z) \equiv \frac{d_L^{\text{GW}}(z)}{d_L^{\text{EM}}(z)}$$



Alice Garoffolo

Detecting DE fluctuations ?

We pick up additional, theory-dependent corrections from inhomogeneities along the line of sight:

$$\frac{\Delta d_L^{\text{SN}}}{\bar{d}_L^{\text{SN}}} = -\kappa - (\phi + \psi) + \frac{1}{\chi} \int_0^\chi d\tilde{\chi} (\phi + \psi) + \phi \left(\frac{1}{\chi \mathcal{H}} \right) + v_{\parallel} \left(1 - \frac{1}{\mathcal{H}\chi} \right) - \left(1 - \frac{1}{\chi \mathcal{H}} \right) \int_0^\chi d\tilde{\chi} (\phi' + \psi')$$



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$$\frac{\Delta d_L^{\text{SN}}}{\bar{d}_L^{\text{SN}}} = -\kappa - (\phi + \psi) + \frac{1}{\chi} \int_0^\chi d\tilde{\chi} (\phi + \psi) + \phi \left(\frac{1}{\chi \mathcal{H}} \right) + v_{\parallel} \left(1 - \frac{1}{\mathcal{H}\chi} \right) - \left(1 - \frac{1}{\chi \mathcal{H}} \right) \int_0^\chi d\tilde{\chi} (\phi' + \psi')$$

$$\frac{\Delta d_L^{\text{GW}}}{\bar{d}_L^{\text{GW}}} = -\kappa - (\phi + \psi) + \frac{1}{\chi} \int_0^\chi d\tilde{\chi} (\phi + \psi) + \phi \left(\frac{1}{\chi \mathcal{H}} - \frac{M'_P}{\mathcal{H}M_P} \right) + v_{\parallel} \left(1 - \frac{1}{\mathcal{H}\chi} + \frac{M'_P}{\mathcal{H}M_P} \right) - \left(1 - \frac{1}{\chi \mathcal{H}} - \frac{M'_P}{\mathcal{H}M_P} \right) \int_0^\chi d\tilde{\chi} (\phi' + \psi') + \frac{M_{P,\varphi}}{M_P} \delta\varphi + \frac{M_{P,X}}{M_P} \delta X$$



Alice Garoffolo

Detecting DE fluctuations ?

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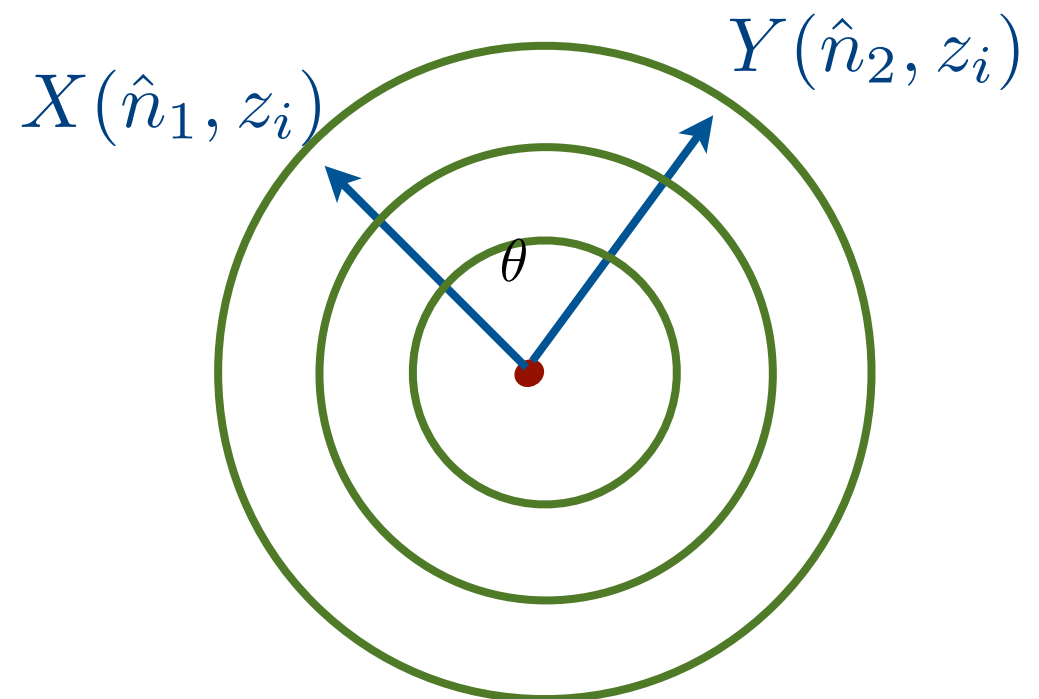
$$\Delta_\varphi(\hat{\theta}, z) \equiv \frac{\Delta d_L^{\text{SN}}(\hat{\theta}, z)}{\bar{d}_L^{\text{SN}}} - \frac{\Delta d_L^{\text{GW}}(\hat{\theta}, z)}{\bar{d}_L^{\text{GW}}} = \frac{M'_P}{\mathcal{H}M_P} \left(\phi - v_{\parallel} + \int_0^\chi d\tilde{\chi} (\phi' + \psi') \right) - \frac{M_{P,\varphi}}{M_P} \delta\varphi - \frac{M_{P,X}}{M_P} \delta X$$

Detecting DE fluctuations ?

$$\langle X Y \rangle_\theta$$

X,Y: GWs, SNIa

$$\downarrow$$
$$C_\ell^{XY}$$



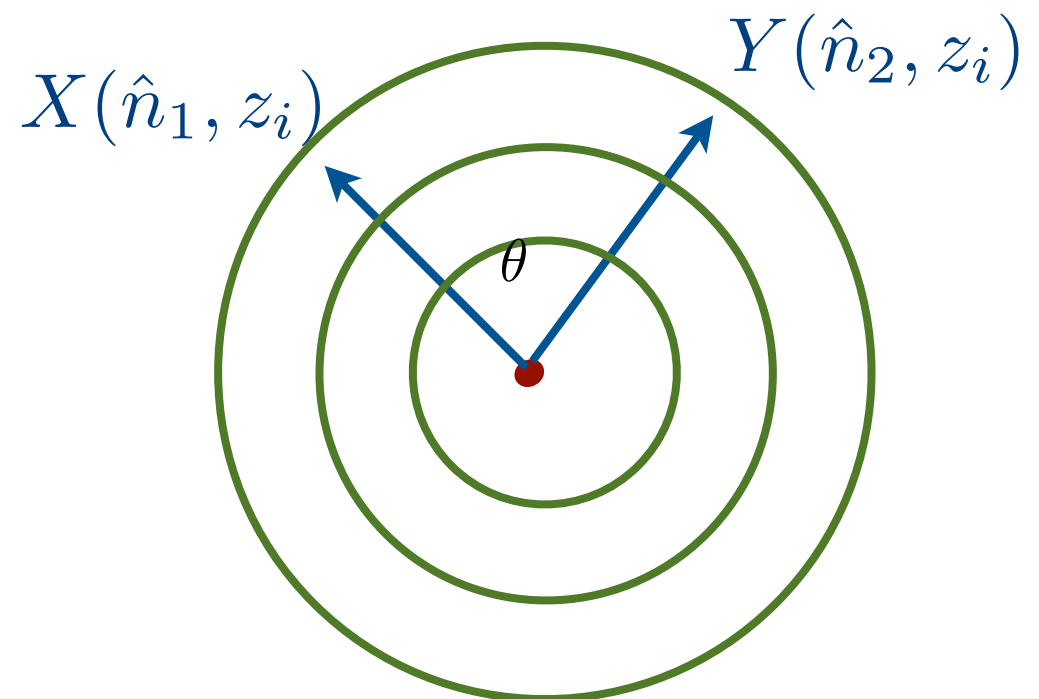
Detecting DE fluctuations ?

$$\langle X Y \rangle_\theta$$

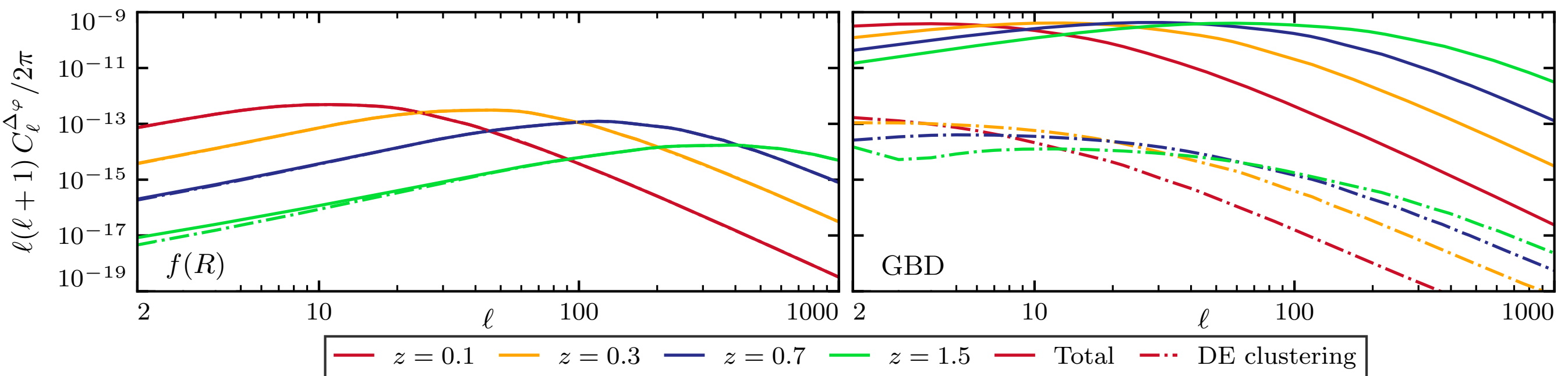
X,Y: GWs, SNIa

$$\downarrow$$

$$C_\ell^{XY}$$



$$C_\ell^{\Delta\varphi} = C_\ell^{\text{SN}} + C_\ell^{\text{GW}} - 2C_\ell^{\text{SN-GW}}$$

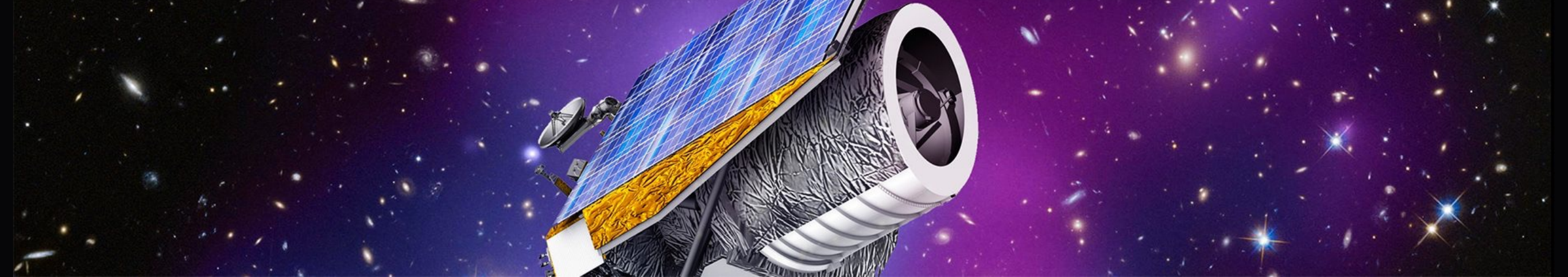


futuristic, but exciting!

$$N_i^{\text{eff}} = \frac{N_i}{\sigma_{d_L}^2 / d_L^2}$$

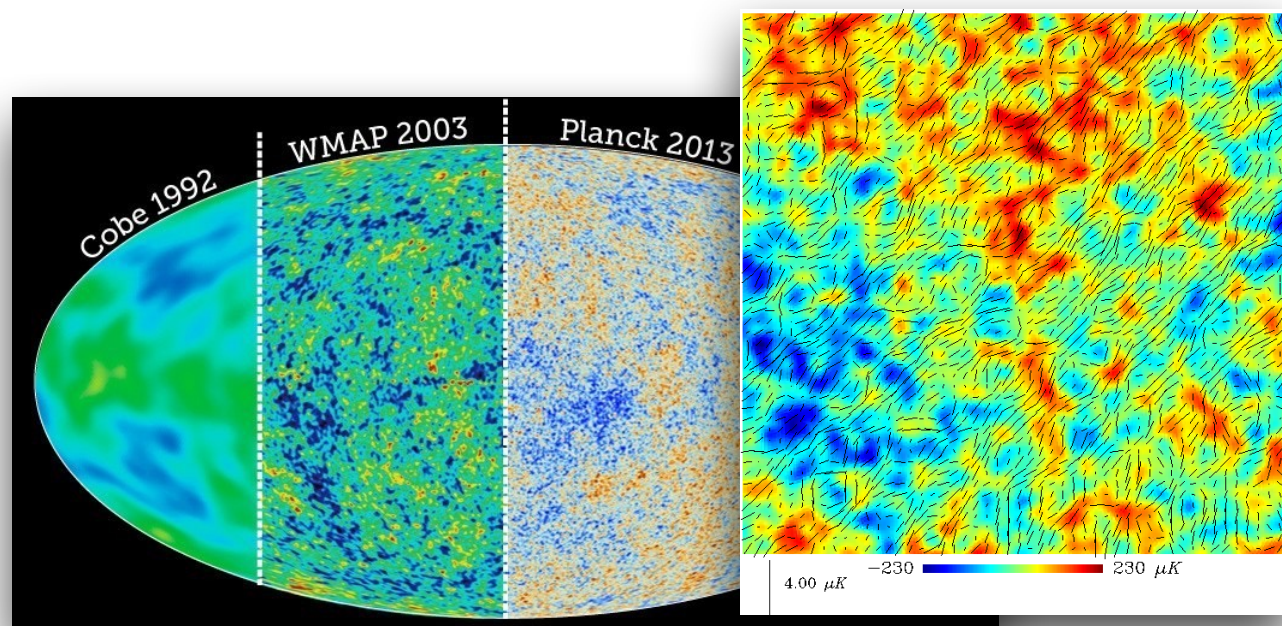
	GR	$f(R)$		GBD	
	$N_{\text{GW}}^{\text{eff}}$	$N_{\text{GW}}^{\text{eff}}$	$N_{\Delta\varphi}^{\text{eff}}$	$N_{\text{GW}}^{\text{eff}}$	$N_{\Delta\varphi}^{\text{eff}}$
$z = 0.1$	10^7	10^7	10^{14}	10^7	10^{12}
$z = 0.3$	10^8	10^8	10^{15}	10^8	10^{11}
$z = 0.7$	10^8	10^8	10^{16}	10^8	10^{12}
$z = 1.5$	10^7	10^7	10^{17}	10^7	10^{12}
w/o z	10^7	10^7	10^{19}	10^7	10^{14}

TABLE I. Effective number of events for a $5\text{-}\sigma$ detection of C_ℓ^{GW} and $C_\ell^{\Delta\varphi}$.

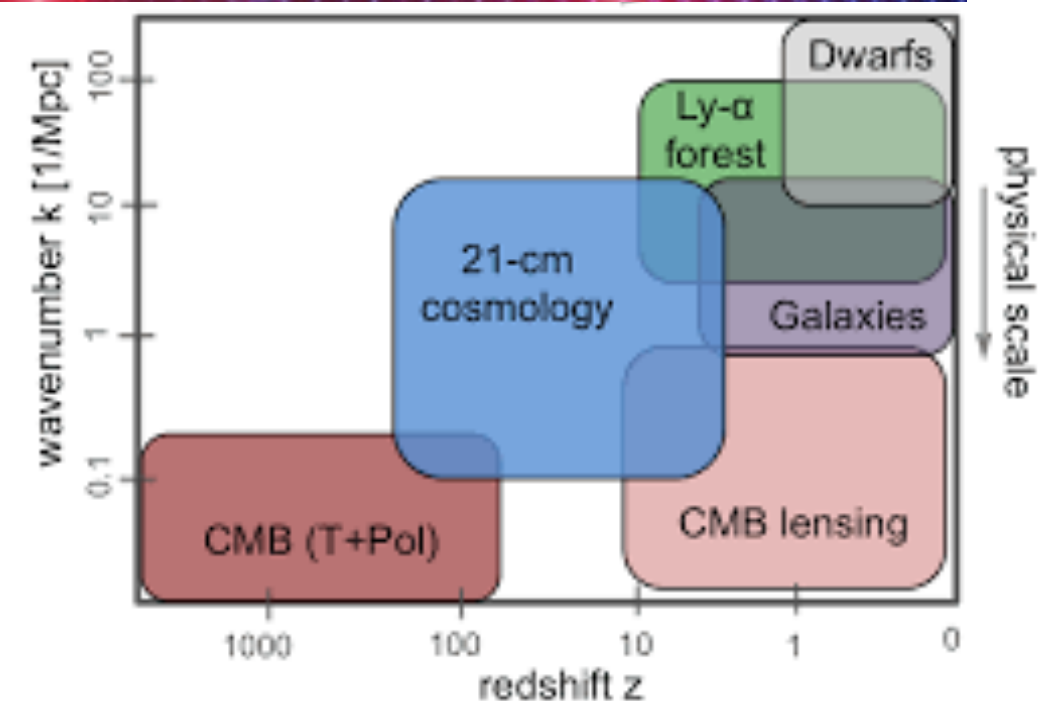


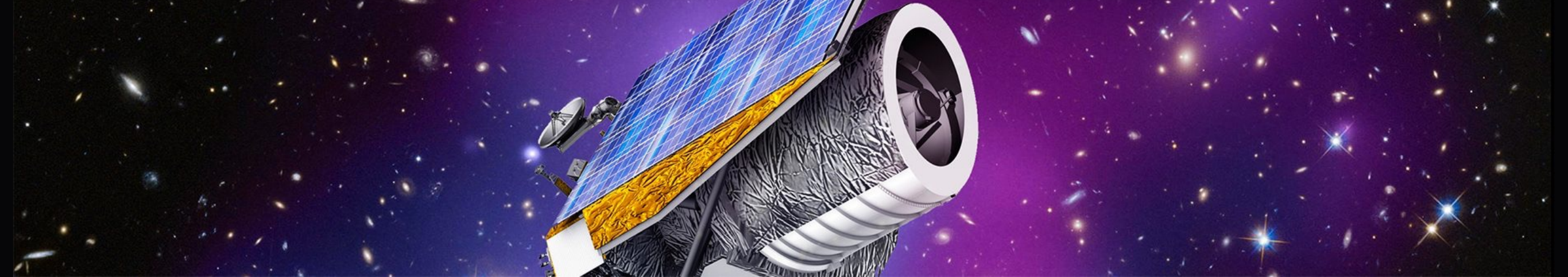
Towards precise and accurate Cosmology

Cosmology is a versatile tool that can test broad classes of theoretical scenarios.
The next decade of observations will see a tremendous leap in sensitivity.
Synergy will be the key!

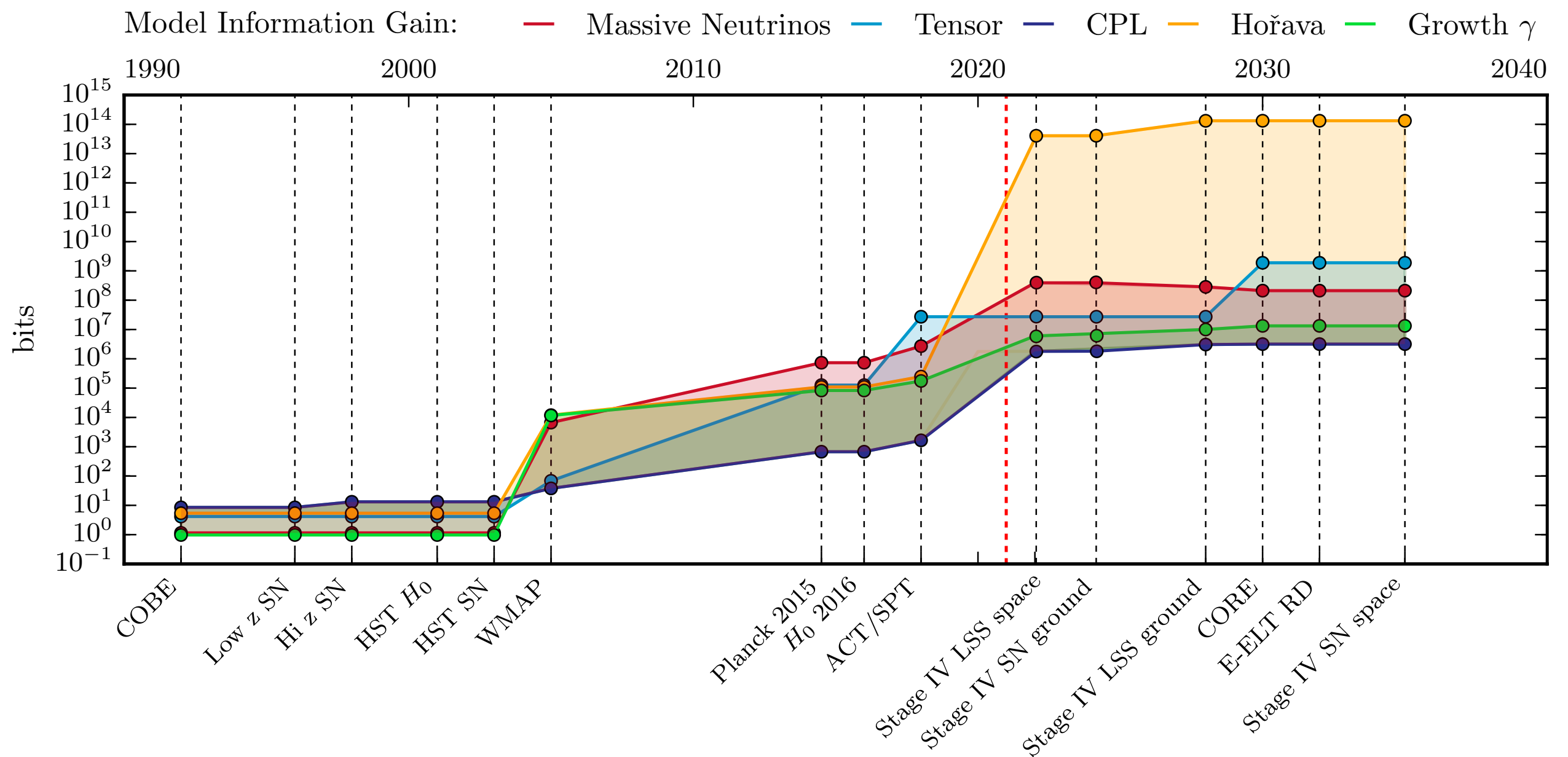


CMB S4

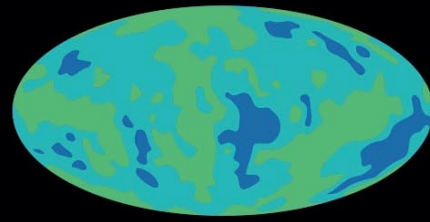




Towards precise and accurate Cosmology



courtesy of M. Raveri, & M. Martinelli (CosmicFish, <https://cosmicfish.github.io/>)



KEEP
CALM
AND
TEST
GRAVITY

the EFTCMB team

THANK YOU
for hosting me !