

GRAVITATIONAL WAVE ASTRONOMY

LECTURE 2: GRAVITATIONAL-WAVE DATA ANALYSIS FOR COMPACT BINARIES

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SUMMARY

SUMMARY LECTURE 1

- Gravitational waves are propagating oscillations of a gravitational field generated by accelerating masses
	- They change the proper separation between freely-falling test bodies
	- GWs carry energy & momentum from the source
- Spacetime is stiff "extreme" events are needed to produce a measurable strain
	- Compact binary mergers, CCSNe, rotating neutron stars, stochastic GW background
- Operating GW detectors are km-scale (sophisticated) Michelson interferometers
	- The sensitivity is characterised by the noise power spectral density (PSD)
	- ▸ Current generation of GW detectors can measure length changes *δL* ≃ 10−18*m*
	- ▸ Observe GWs with frequencies between 20-2000 Hz

DATA ANALYSIS FOR COMPACT BINARIES

- GW detection
- Parameter estimation
- Modelling gravitational waves

(Some) Recommended literature:

★ Luc Blanchet, Living Reviews in Relativity: <https://link.springer.com/article/10.12942/lrr-2014-2>

★ Buonanno & Sathyaprakash: <https://arxiv.org/abs/1410.7832>

★ Talbot & Thrane:<https://arxiv.org/abs/1809.02293>

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SENSITIVITY

DETECTION OF GWS 4

- The sensitivity of GW detector is characterised by the power spectral density (PSD) of its noise background in the absence of a GW signal.
-

[Living Rev. Rel. 2019 22:2]

GW OBSERVATORIES

DETECTOR NETWORK

- ‣ Kilometre-scale interferometres
- ‣ Sensitive to GWs between a few Hz to a few kHz
- ‣ Simultaneous detection increases detection confidence
- ‣ Improved sky localisation & polarisation measurement
- ‣ Increased duty cycle

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THE GW PIPELINE - A VERY SIMPLIFIED SCHEMATIC 6

GW OBSERVATORIES

LOCALISATION

- ▸ Individual GW detectors are omnidirectional: poor localisation! Sensitivity depends on location, polarisation and frequency
- Simultaneously operating observatories allow for triangulation via arrival time differences

[Sathyaprakash&Schutz, LRR]

GW190425

F = sin² ✓*,* (60) $f(90\% \text{ Cl} \sim 8300 \text{ deg}^2)$ $(80, 90)$ with the 9000 confidence region bounded by the thin documentation by the thin documentation by the thin documentation $(10, 10, 0)$ 90% CI ~ 8300 deg2

[LVK Observing Scenarios & LVK discovery papers]

GW SEARCHES

MODELLED SEARCHES

▶ Optimal detection strategy for a modelled signal in stationary Gaussian noise = matched filtering

MMM WM

Ωˆ(*t*) lte

strain data d template waveforms h SNR time series

Optimal filter: $\rho =$ $\langle d|h \rangle$ $\sqrt{\langle h|h \rangle}$ signal-to-noise ratio (SNR)

 $\langle a|b \rangle$

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$$
=4\int_0^\infty df \frac{\tilde{a}(f)\tilde{b}^*(f)}{S_n(f)}
$$

TEMPLATE BANKS *M* = (*h* $\hat{h}(\bar{\theta})|\hat{h}(\bar{\theta}+\Delta\theta))$

GW SEARCHES

- **Large collection of theoretical waveforms templates:**
- ▶ Construction: hybrid method $h_{+}(\tau_{\text{obs}}) = h_{+}(\tau_{\text{obs}})$ $1 + \cos^2 t$ $\frac{1}{2}(\partial \theta) + \frac{\partial}{\partial \theta} \frac{\partial}{\partial \theta} \left[\Phi(\tau_{\text{obs}}) \right],$ $h_{\times}(\tau_{\text{obs}}) = h_c(\tau_{\text{obs}}) \cos \theta_2 \sin \left[\Phi(\tau_{\text{obs}})\right]$ $\frac{d}{dt} \oint_{\mathbf{H}} f(\mathcal{F}_{0}(\mathbf{r}) \hat{\mathbf{r}}) \left(\frac{1}{\theta} \right) \left(\frac{\partial \mathcal{F}}{\partial \mathbf{r}} \right) d\mathbf{r}$ $\frac{\partial}{\partial \theta^i} \mathbb{X}^i_{\theta^i}$ 1 2 $\partial^2 M$ 1 @²*M*
	- **Geometric lattice for known proprighted area)** $=$ dub $\frac{1}{2}$ o ivi ivident
- \blacktriangleright Metric = mismatch between negighbouring $\limsup_{n \to \infty} \frac{1 - \langle \psi_n \theta_i \rangle + \langle \psi_n \theta_j \rangle + \langle \psi_n \$ $\mathbf{F}(t)$ Stochastic placement \blacktriangleright , What's included? ̂ ̂ $f_{\text{sw}}^{(s)} = (1+z) f_{\text{sw}}^{(\text{obs})}$ ismatch between neighbouring $\frac{2}{h} \left(\tau \right)^2 + \frac{\cos^2 k}{\cos^2 k} \frac{J}{\cos^2 (h(\tau))}$ $h_{\text{\textbf{+}}} (t)$ Stochas $a(t)$ ^r ${\mathcal A}$ (acem/er) 4×76 68 2 \bar{t} 2 h_c ($\sqrt{\frac{1}{2\pi}}$ $\frac{1}{2}$ dt' $\vec{f}_{\text{gw}}^{\text{source}}(t-r/c)$ $h_{\times}(t) = \frac{What{h}_{\times}(t)}{M_{\star}}$ *a*(*t*) *r* $\begin{bmatrix} \frac{\partial}{\partial s} & h_c(\tau_{\text{obs}}) \ -\frac{\partial}{\partial s} & f_c(\tau_{\text{obs}}) \end{bmatrix} = h_c(\tau_{\text{obs}})$ $\bar{\partial} \pi$ $\sum_{i=1}^{\infty} \sum_{j=1}^{t} r^{j}$ $\frac{d}{dx} \int_{\text{sw}}^{\text{source}} (t - r/c)$ hat the $\theta_1 \pm \sqrt{2}$ $1 + \frac{3}{2}$ $2\overleftrightarrow{\partial}\theta^2\overleftrightarrow{\partial}\theta^3\overleftrightarrow{\partial}\theta^2$ h_{\times} (τ_{obs}) ϵ h_c (τ_{obs}) ϵ ⁶⁵*l*/ \sin [Φ (τ_{obs})] $f_{\text{gw}}^{(s)} = f_{\text{gw}}^{(s)} \exp\left\{\frac{\mathbf{p}}{2}\right\} \mathbf{f}_{\text{gw}}^{(s)}$ \dot{J} $\overline{2}$ $\partial^2_\bullet M$ $\partial \theta^i \partial \theta^j$ $h_{\text{eff}}(\tau_{\text{obs}}) = h_c(\tau_{\text{obs}})$ $1 + \cos^2 \nu$ $\frac{1}{2} f_{\text{gw}}^{(s)} t^{\frac{1}{2}} \mathcal{N} \mathcal{N} \left(\frac{1}{2} \mathcal{N} \mathcal{N} \mathcal{N} \right),$ $\eta(x)\bar{x}_{\text{obs}}^{\text{A}}\mathcal{V} = \eta_c^{\text{C}}\mathcal{E}_{\text{obs}}^{\text{cos}}\mathcal{V}^{\text{diag}}\mathcal{E}_{\text{OS}}^{\text{diag}}[\mathcal{P}_{\text{S}}^{\text{H}}\mathcal{P}(\Phi(\mathcal{H}_{\text{obs}}^{\text{U}}))]$ \mathcal{H}_{obs} \mathcal{H}_{COS} $\mathcal{H}_{\text{Sinf}}$ \mathcal{H}_{COS} h_{α} $\left(t_{\beta}^{\text{ret}} \right)$
- P Schmidt, Wniv. of Birmingham ▸ Component masses \blacktriangleright Aligned-spins \blacktriangleright Quadr \natural polar mode $p_t = h_c(\tau_{\text{obs}}) - e^2 \frac{g}{2} \longrightarrow \cos\left[\frac{\Phi_t(\tau_{\text{obs}})}{r^t + r\rho_{\text{obs}}}\right]$ b a $\bm{\Theta}^{\mathrm{ret}}$ $\frac{n_+(\iota)}{\gcd(\frac{1}{\sigma_0})} \frac{1}{\gcd(\frac{1}{\sigma_0})}$ *d*L $\frac{A(t) \pm B(t)}{A(t) \pm C(t)}$ e^{2} $\frac{5}{3}$ π 9 (b) $\frac{2}{7}$ π s (b) $\frac{2}{5}$ π) \sum_{c} $h_+(t) = \frac{1}{a\mathbf{f} t s} \mathcal{A}(t, t) + \mathcal{C}^{0.52}_{a\mathbf{f} t s} \left(\frac{1 + \cos^2 t}{\sin^2 t} \right) \frac{1 + \cos^2 t}{\sin^2 t}$ $h_c(t_s^{\rm ret}$ $\mathcal{L}_{\text{ref}}^{\text{left}}(\mathbf{f}_{\text{obs}}) = \frac{1}{\alpha \mathbf{v}^2} + \frac{1}{\alpha \mathbf{v}^2} \sum_{k=1}^{N} \frac{1}{\alpha} \sum_{k=1}^{N} \frac{1}{\alpha} \sum_{k=1}^{N} \frac{1}{\alpha} \left(\frac{1}{\alpha} \sum_{k=1}^{N} \frac{1}{\alpha} \sum_{k=1}^{N} \frac{1}{\alpha} \right)$ $a(f_{\text{emis}}^{\text{L}})$ \mathscr{C} OS *c* $\sqrt{\frac{2^{t/3}}{2}}$ $A(t)$ $\frac{h_c(t) \partial_t \mathbf{y}}{\partial t}$ source *c* $\frac{1}{2} \oint \frac{1}{2} \int \frac{1$ q $\left(\frac{1}{2}\right)$ $A(t, \mathcal{T} | \mathcal{L})$ $1+\cos^2\iota$ 2
2 $q\phi$ sbs $2q$ $\int_0^{t-r/c}$ $h_{\chi}(t) = \frac{1}{2\pi\epsilon}$ *a*(*t*) *r* (τ_{obs}) $\frac{e^2}{2}$ $\left(\frac{t}{\tau} - \frac{r}{c}\right)$ cos ι_{obs} sin (τ_{obs}) $\iint t \rightarrow \rho d$ $h_c(t)$ ret $\frac{2}{6}$ tet $\frac{2}{\sqrt{3}}$ sop $\frac{d}{dt}$ $\frac{1}{4}$ 5*/*3 $G{M_c}$ \int_0^∞ 5*/3*^g $\pi f^{(\rm obs)}$ (*t*) $\frac{df(\xi)}{d\eta}$ $\frac{f(\xi)}{d\eta}$ $\frac{f(\xi)}{d\eta}$ $\frac{f(\xi)}{d\eta}$ $\frac{f(\xi)}{d\eta}$ $h_+(t)$ $a(t)$ $\mathcal{A}(\mathcal{t} - r/\mathcal{e})$ $\frac{1}{t} + \cos^2 t$ \sum $\frac{\sqrt{2}}{60}$ 12π $\int_0^t \frac{f \psi_r}{c}$ $dt_{\rm s}^{\prime}$ $\mathcal{H} \circ \mathcal{H} \circ \mathcal{$ $\frac{1}{h}$ $\begin{pmatrix} f \\ f \end{pmatrix}$ = $\frac{Z_{\text{c}}}{\sqrt{2}}$ cos $\frac{Z_{\text{c}}}{\sqrt{2}}$ *T*^t/ t^l the contract of t $\hat{h}^{\text{top}}_{+}(\tau_{\text{obs}}) = \frac{1}{2}\hat{h}^{\text{in}}_{c}(\tau_{\text{obs}})$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{2}{2} \int_{\cos t}^{\cos t} \cos \left[\frac{\Phi}{f}(\tau_{\text{p}}) \right]$, *f*(*s*)

 ϵ^2

 $\Omega +$

 $\hat{h}(\bar{\theta}) = h(\bar{\theta}) / \sqrt{(h(\bar{\theta}))h(\bar{\theta})}$ $\overline{}$ $(h(\bar{\theta})|h(\bar{\theta}))$

◆²*/*³

GW SEARCHES

EVENT SELECTION

- Apply any data excursions
- Filter each data stream individually
- Apply threshold SNR cut and cluster to generate triggers
- **•** Perform χ^2 -test to check signal consistency
- Check time & parameter coincidence between different IFOs
- Apply data quality vetoes
- Surviving triggers form GW candidates
- ▸ Use time shifts to calculate the false-alarm-rate (FAR) of coincident triggers. Resulting triggers form the background, which is used to determine the significance of the foreground candidates. int_{right} $\sqrt{\Gamma(\text{AD})}$ of

Distribution of SNR contributions by frequency band Example: SNR accumulation as signal-consistency test

GW SEARCHES

SIGNIFICANCE

- **FAR** (caution: varies!): e.g. O2 1/30 days (expect \sim 2 false alarms per month)
- \blacktriangleright **Astrophysical probability**: assumes 4 categories A_i (terrestrial, BBH, BNS, NSBH) each described by a Poisson process (inhomogeneous Poisson mixture model) with mean Λ_i [Farr+, 2013]
	- **•** Ranking statistic distribution: $p(x | A_i, \{\theta\})$
	- $\Lambda_i = \langle VT \rangle_{\{\theta\}} \times R_i$
		- Observable spacetime volume (selection effects!)
	- ▸ Models are marginalised over the counts with the ranking statistic distribution fixed at the value of the ranking statistic of the candidate

[LVK, GWTC-1] HV GMTC-11

while also probing noise triggers. The contract of the contract of the contract of the contract of the contract of

triggers with lower \sim 11

UNMODELLED SEARCHES

- Search that identifies coincident excess power in the time-frequency representation of the strain data
	- Identifies events that are coherent in multiple detectors and reconstructs the source sky location and signal waveforms by using the constrained maximum likelihood method
	- Does not rely on waveform models
	- Sensitive to a wide range of short-duration transient signals ("bursts")
	- Weak assumption of "chirpyness" of the signal
- Detection statistic: coherent energy constructed via cross-correlation $E_c \propto \rho_c$

GW SEARCHES 12

[Salemi+, 2019]

PARAMETER ESTIMATION \blacksquare **PARAMETER ESTIMATION** regime and confirm predictions of \mathbb{R}^n **HANGELIST CONSERVATION**

regime and confirm predictions of general relativity for the

SOURCE CHARACTERISATION $t \sim t$ and $\sim t$ shown in Fig. 1. The initial GW150914 shown in Fig. **PRILIDAE PHADAPTEDICATION** nonlinear dynamics of highly disturbed black holes. The dynamics of highly disturbed black holes. The dynamics IVL VIIF gravitational-wave transients \mathcal{A}^{max} and was reported with \mathcal{A}^{max} \blacksquare matched-filter analyses that use relativistic models of compart binary waveforms $\mathbf{4}$

Calibratad (rawich) ctrain data t thar minutes of data acquisition \mathcal{A} Calibrated (raw-ish) strain data

filters to remove the strong instrumental spectral lines seen in the Fig. 3 spectra. Top row, left: H1 strain. Top row, right: L1 strain.

give our best estimate for the parameters describing the parameters describing the parameters describing the p

In Table I, we also indicate the sensitive our results are sure to the sensitive our results are sure to the sensitive our results are sensitive our results and the sensitive our results are sensitive our results and the s

filters to remove the strong instrumental spectral lines seen in the Fig. 3 spectral lines seen in the Fig. 3 spectral lines seen in the Fig. 3 spectral lines seen in the \mathcal{H}

and management

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[LVC, GW150914 discovery papers] which produces the sharp cut in the sharp cut in the two-dimensional distribution. The two-dimensional distrib
The two-dimensional distribution of two-dimensional distribution. The two-dimensional distribution of two-dime

−4M⊙, as shown in Table I and Fig. 1.1 as shown in Table I and Fig. 1.1 as shown in Table I and Fig. 1.1 as sh
1.1 as shown in Table I and Fig. 1.1 as shown in Table I and Fig. 1.1 as shown in Table I and Fig. 1.1 as show

GW DATA ANALYSIS

MEASURING BINARY PARAMETERS

▸ GW signal encodes the astrophysical information: masses, spins, tides, location, orientation of the orbit

▸ Use Bayesian inference to infer the source parameters *θ*:

Posterior	$p(\theta \mid d) = \frac{\mathcal{L}(d \mid \theta)\pi(\theta)}{\mathcal{Z}}$
Probability	$p(\theta \mid d) = \frac{\mathcal{L}(d \mid \theta)\pi(\theta)}{\mathcal{Z}}$

 $R(h) = F_+(\alpha, \delta, \psi)h_+ + F_\times(\alpha, \delta, \psi)h_\times$

PARAMETER ESTIMATION

THE DATA

For an incident signal, a GW detector records the following data:

Real-valued detector response to a GW:

THE NOISE

- Three key assumptions:
	- Gaussian with zero mean and known variance

 $p(n) \propto \exp \left(-\frac{1}{2}\right)$

 $C_{ij} =$

Independent between frequency bins and detectors

 $S_n(f_i)\delta_{ij}$ power spectral density

1

2

noise correlation matrix

LIKELIHOOD

- We are free to choose the likelihood
- variance:

 $-\frac{2|\tilde{h}_i(\theta) - \tilde{d}_i|}{TS_{i}(f_i)}$ $\frac{d_i(\theta) - \tilde{d}_i|^2}{TS_n(f_i)} - \frac{1}{2}\log\left(\pi TS_n(f_i)/2\right)$! waveform template

$$
\mathcal{L}(d|\theta,H,S_n(f))=\exp\left(\sum_i -
$$

- Likelihood compares theoretical models against the data
- For a network of GW detectors:

$$
\mathcal{L}(d_{NW}|\theta, H, S_{nNW}(f)) = \prod_{i \in NW} \mathcal{L}(d_i|\theta, H, S_{ni}(f))
$$

▶ In GW astronomy we assume a likelihood associated with stationary Gaussian noise that has zero mean and a known

STATIONARY & GAUSSIAN?

- Data cleaning
-

DETECTION OF GWS 18

Some example "glitches"

25

 -20

L5

Normalized energy

<https://www.zooniverse.org/projects/zooniverse/gravity-spy>

[Macas+, 2022] ρ Ω_{Z} legalization of a CW150 α strain diata (a) show in the glitch in the time domain. Treature to the bilp glitch central 137 $deg²$. (a) Str localization of a $\text{CN}/150014$ like event injected at $t + 20$ ms α is not improve the nondetector the only gradient central time t_0 . The gravitatione area is 127 deg^2 $\frac{157}{\text{deg}}$. $\frac{1}{20}$ and $\frac{1}{2}$ H₁ $\frac{1}{20}$ we have the theorem 3 kHz so that the theorem 3 kH $\left[\text{M}_{2c25} + 2022\right]$

not visible in the bottom parameters in the bottom parameters \mathbf{r}

GAUSSIANITY Additionally, a short instrumental noise transient

- Real detector noise often contains non-Gaussianities ("glitches") $T_{\rm eff}$ transient noise, or glitch \sim glitch \sim glitch \sim glitch \sim glitch and \sim glitch and \sim annues (ginches) saturation in the digital-to-analog converter to-analog converter
	- ▸ O3: 24% of GW candidates near glitches Ω z4% of GW candidates hear glitches
	- ▸ We can modify the likelihood or remove them from the data $\mathbf{u} \cdot \mathbf{c} = \mathbf{u} \cdot \mathbf{u}$ and the results presented in the result can modify the likelinood or remove $x \text{ from the data}$ treatment of other high-amplitude glitches used in the second in the
- ▸ Prominent example: GW170817 ont ovemplo: $G11/170817$ subtraction description description description description of the glitch was modeled above, the glitch was modeled
- ▸ Noise transients impact the measurement: s Following the procedure for prior gravitation.
	- ▸ Systematic errors in sky location environmental disturbance observed by LIGO environmenemalic enors in sky location
	- ▸ Systematic biases in source parameters The Virgo data, used for sky localization and an ematic plases in source parameters,

 $Fvamb: G11/170917ovcalian$ Example: GW170817 overlapped with a blip glitch in LIGO-Livingston are shown relative to August 17, 2017 12∶41:04 UTC. Top panel: top panel: top panel: top panel: top panel: top
Distribution of the UTC. Top panel: top pane

STATIONARITY

- The noise floor in GW detectors varies over time
	- \blacktriangleright $S_n(f)$ changes between events
	- Compute the PSD for every event to take PSD drift into account

Note: Uncertainty in the PSD estimation can also be taken into account via marginalisation by introducing extra parameters into the noise model

Credit: K. Chatziioannou

Hanford

CALIBRATION

- Calibration quantifies the detector's response to incident GWs
	- Miscalibration results in biased strain data!
- Parameterised model for calibration uncertainty
	- Marginalised over during parameter estimation
- Calibration uncertainties are not the limiting factor in current GW observations
	- Statistical uncertainty from detector noise dominates

[Cahillane+, 2017]]

Credit: LIGO/Virgo

PARAMETER ESTIMATION

THE SIGNAL MODEL

- \triangleright Need a signal model $h(\theta)$ to compare against the data when evaluating the likelihood
- Circular binary black hole (15D):

- Eccentric orbits: requires two additional parameters to describe the ellipse and its orientation
- Binary neutron star: θ_{BBH} + parameters that characterise the tidal response of the star

intrinsic

$$
\theta_{\rm BBH} = \{m_1, m_2, \vec{\chi}_1, \vec{\chi}_2, D_L, \iota, \alpha, \delta, \psi, \phi_c, t_c\}
$$

extrinsic

THE PRIOR

- Uninformative vs. astrophysically motivated priors
	- Population priors
- "Standard" GW priors:
	- ▸ Mass priors: often sample uniform in chirp mass ${\mathscr M}_c = \eta^{3/5} M$ & symmetric mass ratio $\eta = (m_1 m_2)/M$
	- ▶ Distance: $p(D_L) \propto D_L^2$ (uniform in "luminosity" volume)
	- Spins magnitudes: uniform
	- Spin orientation: isotropic (i.e. uniform in cos(tilt)

SPYCBC

TRY IT YOURSELF :-)

- All LVK analysis software is publicly available: <https://git.ligo.org/lscsoft/lalsuite>
- Tutorials available from the Gravitational Wave Open Science Centre:<https://gwosc.org/tutorials/>
- Perform a matched filter search: <https://pycbc.org/>
- User-friendly Python inference package Bilby:
	- Source code:<https://git.ligo.org/lscsoft/bilby>
	- Documentation: <https://lscsoft.docs.ligo.org/bilby/>

WAVEFORMS

GRAVITATIONAL WAVES FROM COMPACT BINARIES

- ▶ Signal "sweeps" through the detector's sensitivity band
- ▸ Depending on the total mass of the binary, the merger regime is visible
	- ▸ Inspiral-merger-ringdown (IMR) waveforms are key
	- ▸ Need accurate theoretical model to infer the science

 $f[Hz]$

"chirp" - encodes the fundamental source properties

WAVEFORMS

precessing spins

MMMMMMMMMMMMMMMMM

AND MUNICIPAL POSET SCHEMINGERS

orbital eccentricity

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MODELLING TECHNIQUES VAAAAANA

WAVEFORMS

Courtesy: H. Pfeiffer

MODELLING TECHNIQUES

ANALYTICAL RELATIVITY TOOLS

- Post-Newtonian (PN) theory
	- Weak fields & slow motion

- Post-Minkowskian (PM) theory
	- ▸ Weak fields, arbitrary velocities

▸ (uncalibrated) effective-one-body (EOB)

Effective field theory (EFT)

- ▶ Scattering amplitudes
	- ▸ Hyperbolic motion
	- ▶ Map to the bound case

- ▸ Gravitational self-force (GSF)
	- ▸ Very unequal mass ratios

- ▸ (Classical) Perturbation theory
	- ▸ Ringdown

"Tutti Frutti"

MODELLING TECHNIQUES

MODELLING THE INSPIRAL

▸ Multi-length-scale problem

Near zone: Equations of motion Post-Newtonian Weak fields & slow motion Source multipoles

from Buonanno&Sathyaprakash *from Buonanno&Sathyaprakash*

[See Blanchet's LRR review]

GW DATA ANALYSIS

GRAVITATIONAL WAVEFORMS

A gravitational-wave signal can be written as:

 $\overline{1}$

The GW phase encodes the physics:

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Spin effects at 1.5PN Slide credit: G. Pratten

- ▸ Only few analytic solutions to Einstein field equations known
	- No analytic solutions to the general-relativistic two-body problem
	- ▸ NR is a key ingredient to understanding GW observations!
- ▸ Breakthrough in 2005 *[Pretorius, Baker+, Campanelli+]*:
	- ▸ First stable binary evolutions
	- Extraction of the GW signal
	- Computationally expensive
	- Time consuming
- Many challenges remain!

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NUMERICAL RELATIVITY (NR)

resolution (due both to the mesh refinement structure to the mesh refinement structure of the mesh refinement s

MODELLING TECHNIQUES

[[]Boyle+ inc. PS, CQG]