



GRAVITATIONAL WAVE ASTRONOMY

# LECTURE 2: GRAVITATIONAL-WAVE DATA ANALYSIS FOR COMPACT BINARIES

PATRICIA SCHMIDT ISCRA 2024





#### SUMMARY

## **SUMMARY LECTURE 1**

- Gravitational waves are propagating oscillations of a gravitational field generated by accelerating masses
  - They change the proper separation between freely-falling test bodies
  - GWs carry energy & momentum from the source
- Spacetime is stiff "extreme" events are needed to produce a measurable strain
  - Compact binary mergers, CCSNe, rotating neutron stars, stochastic GW background
- Operating GW detectors are km-scale (sophisticated) Michelson interferometers
  - The **sensitivity** is characterised by the **noise power spectral density (PSD)**
  - Current generation of GW detectors can measure length changes  $\delta L \simeq 10^{-18} m$
  - Observe GWs with frequencies between 20-2000 Hz





## DATA ANALYSIS FOR COMPACT BINARIES

(Some) Recommended literature:

**t** Luc Blanchet, Living Reviews in Relativity: <u>https://link.springer.com/article/10.12942/lrr-2014-2</u>

**Buonanno & Sathyaprakash:** <u>https://arxiv.org/abs/1410.7832</u>

**Talbot & Thrane:** <u>https://arxiv.org/abs/1809.02293</u>



Schmidt, Univ. of Birmingham



- **GW** detection
- **Parameter estimation**
- Modelling gravitational waves



- background in the absence of a GW signal.



## **GW OBSERVATORIES**

# **DETECTOR NETWORK**



- Kilometre-scale interferometres
- Sensitive to GWs between a few Hz to a few kHz
- Simultaneous detection increases detection confidence
- Improved sky localisation & polarisation measurement
- Increased duty cycle





## THE GW PIPELINE - A VERY SIMPLIFIED SCHEMATIC





P Schmidt, Univ. of Birmingham



## **GW OBSERVATORIES**

## LOCALISATION

- Individual GW detectors are omnidirectional: poor localisation! Sensitivity depends on location, polarisation and frequency
- Simultaneously operating observatories allow for triangulation via arrival time differences













GW190425

90% CI ~ 8300 deg<sup>2</sup>

| .8 1 |  |  |
|------|--|--|
|      |  |  |
|      |  |  |
|      |  |  |
|      |  |  |
|      |  |  |

# **DELLED SEARCHES**

Optimal detection strategy for a **modelled signal** in stationary Gaussian noise = **matched filtering** 

strain data d

 $\langle d|h \rangle$ **Optimal filter:** signal-to-noise ratio (SNR)



Schmidt, Univ. of Birmingham

template waveforms h

SNR time series



$$\langle a|b\rangle = 4 \int_0^\infty df \frac{\tilde{a}(f)\tilde{b}^*(f)}{S_n(f)}$$



## **TEMPLATE BANK** $= (\hat{h}(\bar{\theta})|\hat{h}(\bar{\theta} + \Delta\theta))$

- Large collection of theoretical waveforms templates:  $\begin{pmatrix} S \\ N \end{pmatrix}_{i} \simeq \frac{(h_{i}|s)}{\sqrt{(h_{i}|h_{i})}}$ Construction: hybrid method  $\hat{h}(\theta) = h_{c}(\tau_{0bs}) \frac{1 + \cos^{2} \iota}{(\theta) + \theta} \frac{\partial ds}{\partial \theta^{i}} \frac{[\Phi(\tau_{0bs})]}{\Delta \theta^{i} \Delta \theta^{i}} \frac{\Delta \theta^{i} \Delta \theta^{j}}{\Delta \theta^{i} \partial \theta^{j}} \Delta \theta^{i} \Delta \theta^{j} \int_{-\infty}^{\infty} df \tilde{A}^{*}(f) \tilde{B}(f) S_{n}^{2}(f)^{\bullet}$ 
  - Geometric lattice for  $\overline{kn}$   $\overline{m}$   $\overline{m}$
- $\begin{array}{l} \text{Metric} = \text{mismatch}^{f^{(s)}}_{\text{pw}} = (1+z) f^{(\text{obs})}_{\text{swhbouring}} \\ \text{templates } 1 \langle \hat{p}_{k} \varphi \not \downarrow \hat{h} \langle \varphi & \downarrow \\ 2 \partial \theta^{i} \partial \theta^{j} & \downarrow \\ 2 \partial \theta^{i} \partial \theta^{j} & \downarrow \\ 2 \partial \theta^{i} \partial \theta^{j} & \downarrow \\ \end{array} \right. \begin{array}{l} 1 + \cos^{2} \mu e^{i \varphi^{\text{ret}}}_{\text{templates}} & 4 \\ \frac{1}{\tau_{\text{obs}}} & \frac{1}{\tau_{\text{obs}}} + \cos^{2} \mu e^{i \varphi^{\text{ret}}}_{\text{templates}} & 4 \\ \frac{1}{\tau_{\text{obs}}} & \frac{1}{\tau_{\text{obs}}} + \cos^{2} \mu e^{i \varphi^{\text{ret}}}_{\text{templates}} & \frac{4}{\tau_{\text{obs}}} & \left(\frac{G \mathcal{H}_{c}}{\tau_{\text{obs}}} + \frac{5/3}{|x_{1}|^{2} + \frac{1}{\tau_{\text{obs}}}} \right)^{2/3}_{\text{templates}} \\ \frac{1}{\tau_{\text{obs}}} & \frac{1}{\tau_{\text{obs}}} + \frac{1}{\tau_{\text{obs}}$  $h_{\star}(t) \operatorname{Stochastic} \operatorname{placem/ent}^{2} \frac{h_{\star}(\tau_{60s})^{2}}{2} = h_{c}(\tau_{obs}) \operatorname{placem/ent}^{2} \frac{h_{\star}(\tau_{60s})^{2}}{2\pi} \int_{0}^{0} \operatorname{placem/ent}^{2} \frac{h_{\star}(\tau_{60s})}{2\pi} \int_{0}^{0} \operatorname{placem/ent}^{2} \frac{h_{\star}(\tau_{60s})^{2}}{2\pi} \int_{0}^{0} \operatorname{placem/ent}^{2} \frac{h_{\star}(\tau_{60s})}{2\pi} \int_{0}^{0} \operatorname{p$

Component masses
$$h_{+}(t) = \frac{1}{a(t) r} \frac{A(t - r/c) \frac{1 + \cos^{2} t}{f(t) + \frac{1}{2} + \frac$$

 $\hat{h}(\bar{\theta}) = h(\bar{\theta}) / \sqrt{(h(\bar{\theta})|h(\bar{\theta}))}$ 





## **EVENT SELECTION**

- Apply any data excursions
- Filter each data stream individually
- Apply threshold SNR cut and cluster to generate triggers
- Perform  $\chi^2$ -test to check signal consistency
- Check time & parameter coincidence between different IFOs
- Apply data quality vetoes
- Surviving triggers form GW candidates
- Use time shifts to calculate the false-alarm-rate (FAR) of coincident triggers. Resulting triggers form the background, which is used to determine the significance of the foreground candidates.





Example: SNR accumulation as signal-consistency test



## SIGNIFICANCE

- FAR (caution: varies!): e.g. O2 1/30 days (expect ~2 false alarms per month)
- **Astrophysical probability**: assumes 4 categories  $A_i$  (terrestrial, BBH, BNS, NSBH) each described by a Poisson process (inhomogeneous Poisson mixture model) with mean  $\Lambda_i$  [Farr+, 2013]
  - Ranking statistic distribution:  $p(x | A_i, \{\theta\})$
  - $\land \quad \Lambda_i = \langle VT \rangle_{\{\theta\}} \times R_i$ 
    - Observable spacetime volume (selection effects!)
  - Models are marginalised over the counts with the ranking statistic distribution fixed at the value of the ranking statistic of the candidate







[LVK, GWTC-1]

## **UNMODELLED SEARCHES**

- Search that identifies coincident excess power in the time-frequency representation of the strain data
  - Identifies events that are coherent in multiple detectors and reconstructs the source sky location and signal waveforms by using the constrained maximum likelihood method
  - Does not rely on waveform models
  - Sensitive to a wide range of short-duration transient signals ("bursts")
  - Weak assumption of "chirpyness" of the signal
- > Detection statistic: coherent energy constructed via cross-correlation  $E_c \propto \rho_c$









[Salemi+, 2019]



## **SOURCE CHARACTERISATION**

#### Calibrated (raw-ish) strain data



3





20

25

30

[LVC, GW150914 discovery papers]

40

35

 $m_1^{
m source}/{
m M}_{\odot}$ 



50

45

#### **GW DATA ANALYSIS**

## **MEASURING BINARY PARAMETERS**

GW signal encodes the astrophysical information: masses, spins, tides, location, orientation of the orbit

Use **Bayesian inference** to infer the source parameters  $\theta$ :

Posterior 
$$p(\theta \mid d) = \frac{\mathcal{L}(d \mid \theta)\pi(d)}{\mathcal{Z}}$$
Evidence







## THE DATA

For an incident signal, a GW detector records the following data:





Real-valued detector response to a GW: 





 $R(h) = F_{+}(\alpha, \delta, \psi)h_{+} + F_{\times}(\alpha, \delta, \psi)h_{\times}$ 



## THE NOISE

- Three key assumptions:
  - Gaussian with zero mean and known variance

 $p(n) \propto \exp\left(-\frac{1}{2}\sum_{i,j}n_i C_{ij}^{-1}n_j\right)$ 



**Independent** between frequency bins and detectors 





noise correlation matrix

 $C_{ij} = \frac{1}{2} S_n(f_i) \delta_{ij}$ power spectral density



## IKFI IHNNN

- We are free to choose the likelihood
- variance:

$$\mathcal{L}(d|\theta, H, S_n(f)) = \exp\left(\sum_{i} -\frac{1}{i}\right)$$

- Likelihood compares theoretical models against the data
- For a network of GW detectors:

$$\mathcal{L}(d_{\mathrm{NW}}|\theta, H, S_{n\mathrm{NW}}(f)) = \prod_{i \in \mathrm{NW}} \mathcal{L}(d_i|\theta, H, S_{ni}(f))$$



In GW astronomy we assume a likelihood associated with stationary Gaussian noise that has zero mean and a known

waveform template  $-\frac{2|\tilde{h}_i(\theta) - \tilde{d}_i|^2}{TS_n(f_i)} - \frac{1}{2}\log\left(\pi TS_n(f_i)/2\right)\right)$ 



## **DETECTION OF GWS**

## **STATIONARY & GAUSSIAN?**

- Data cleaning



https://www.zooniverse.org/projects/zooniverse/gravity-spy



Some example "glitches"







15



## GAUSSIANITY

- Real detector noise often contains **non-Gaussianities** ("glitches")
  - O3: 24% of GW candidates near glitches
  - We can modify the likelihood or remove them from the data
- Prominent example: GW170817
- Noise transients impact the measurement:
  - Systematic errors in sky location
  - Systematic biases in source parameters



500

100

Frequency (Hz)



Example: GW170817 overlapped with a blip glitch in LIGO-Livingston



(a) Sky localisation of a GW150914-like event injected at  $t_0 + 30 \text{ ms}$ relative to the blip glitch central time  $t_0$ . The 90% credible area is  $137 \, deg^2$ . [Macas+, 2022]

## **STATIONARITY**

- The noise floor in GW detectors varies over time
  - $S_n(f)$  changes between events
  - Compute the PSD for every event to take PSD drift into account

Note: Uncertainty in the PSD estimation can also be taken into account via marginalisation by introducing extra parameters into the noise model





Credit: K. Chatziioannou



## CALIBRATION

- Calibration quantifies the detector's response to incident GWs
  - Miscalibration results in biased strain data!
- Parameterised model for calibration uncertainty
  - Marginalised over during parameter estimation
- Calibration uncertainties are not the limiting factor in current GW observations
  - Statistical uncertainty from detector noise dominates





[Cahillane+, 2017]]





# **THE SIGNAL MODEL**

- Need a signal model  $h(\theta)$  to compare against the data when evaluating the likelihood
- Circular binary black hole (15D):

intrinsic

extrinsic

- Eccentric orbits: requires two additional parameters to describe the ellipse and its orientation
- Binary neutron star:  $\theta_{\rm BBH}$  + parameters that characterise the tidal response of the star







Credit: LIGO/Virgo







## THE PRIOR

- Uninformative vs. astrophysically motivated priors
  - Population priors
- "Standard" GW priors:
  - Mass priors: often sample uniform in chirp mass  $\mathcal{M}_c = \eta^{3/5} M$  & symmetric mass ratio  $\eta = (m_1 m_2)/M$
  - > Distance:  $p(D_L) \propto D_L^2$  (uniform in "luminosity" volume)
  - Spins magnitudes: uniform
  - Spin orientation: isotropic (i.e. uniform in cos(tilt)





[LVC, GWTC-1]



## **MEASUREMENT EXAMPLES**



- Chirp mass vs total mass
- Distance-inclination degeneracy
- Individual spins vs effective spins

Impact of SNR (statistical uncertainty)



600

400

0



![](_page_23_Picture_11.jpeg)

# TRY IT YOURSELF :-)

DIY

- All LVK analysis software is publicly available: https://git.ligo.org/lscsoft/lalsuite
- Tutorials available from the Gravitational Wave Open Science Centre: <u>https://gwosc.org/tutorials/</u>
- Perform a matched filter search: <u>https://pycbc.org/</u>
- User-friendly Python inference package Bilby:
  - Source code: <u>https://git.ligo.org/lscsoft/bilby</u>
  - Documentation: <a href="https://lscsoft.docs.ligo.org/bilby/">https://lscsoft.docs.ligo.org/bilby/</a>

![](_page_24_Picture_7.jpeg)

# **OPYCBC**

![](_page_24_Picture_11.jpeg)

![](_page_24_Picture_12.jpeg)

![](_page_24_Picture_13.jpeg)

#### WAVEFORMS

# **GRAVITATIONAL WAVES FROM COMPACT BINARIES**

- Signal "sweeps" through the detector's sensitivity band
- Depending on the total mass of the binary, the merger regime is visible
  - Inspiral-merger-ringdown (IMR) waveforms are key
  - Need accurate theoretical model to infer the science

![](_page_25_Figure_6.jpeg)

![](_page_25_Picture_8.jpeg)

![](_page_25_Figure_10.jpeg)

f[Hz]

![](_page_25_Picture_13.jpeg)

#### WAVEFORMS

# PHENOMENOLOGY

#### no spins or aligned spins

![](_page_26_Figure_3.jpeg)

![](_page_26_Picture_4.jpeg)

![](_page_26_Figure_6.jpeg)

![](_page_26_Picture_8.jpeg)

![](_page_26_Picture_10.jpeg)

![](_page_26_Picture_11.jpeg)

![](_page_26_Picture_12.jpeg)

![](_page_26_Picture_13.jpeg)

#### WAVEFORMS

# MODELLING TECHNIQUES

![](_page_27_Picture_2.jpeg)

![](_page_27_Picture_3.jpeg)

P Schmidt, Univ. of Birmingham

Courtesy: H. Pfeiffer

![](_page_27_Picture_6.jpeg)

## MODELLING TECHNIQUES

## ANALYTICAL RELATIVITY TOOLS

- Post-Newtonian (PN) theory
  - Weak fields & slow motion

- Post-Minkowskian (PM) theory
  - Weak fields, arbitrary velocities

(uncalibrated) **effective-one-body** (EOB)

Effective field theory (EFT)

![](_page_28_Picture_8.jpeg)

- Scattering amplitudes
  - Hyperbolic motion
  - Map to the bound case

- Gravitational self-force (GSF)
  - Very unequal mass ratios

- (Classical) Perturbation theory
  - Ringdown

#### "Tulti Frutti"

![](_page_28_Picture_18.jpeg)

![](_page_28_Picture_19.jpeg)

## **MODELLING TECHNIQUES**

## **MODELLING THE INSPIRAL**

Multi-length-scale problem

Near zone: Equations of motion Post-Newtonian Weak fields & slow motion Source multipoles

![](_page_29_Picture_4.jpeg)

from Buonanno&Sathyaprakash

![](_page_29_Picture_6.jpeg)

![](_page_29_Picture_7.jpeg)

#### [See Blanchet's LRR review]

diagram not to scale

![](_page_29_Picture_11.jpeg)

![](_page_29_Picture_13.jpeg)

#### **GW DATA ANALYSIS**

A gravitational-wave signal can be written as:

![](_page_30_Picture_3.jpeg)

The GW phase encodes the physics:

![](_page_30_Figure_5.jpeg)

![](_page_30_Picture_6.jpeg)

Schmidt, Univ. of Birmingham

Spin effects at 1.5PN

Slide credit: G. Pratten

![](_page_30_Picture_10.jpeg)

## **MODELLING TECHNIQUES**

# NUMERICAL RELATIVITY (NR)

- Only few analytic solutions to Einstein field equations known
  - No analytic solutions to the general-relativistic two-body problem
  - NR is a key ingredient to understanding GW observations!
- Breakthrough in 2005 [Pretorius, Baker+, Campanelli+]:
  - First stable binary evolutions
  - Extraction of the GW signal
  - Computationally expensive
  - Time consuming
- Many challenges remain!

![](_page_31_Picture_11.jpeg)

Schmidt, Univ. of Birmingham

- - Extract gravitational waves

エ Φ  $\operatorname{Re}(\Psi_4) \cdot \Gamma,$ 

![](_page_31_Figure_16.jpeg)

<sup>1/</sup>q[Boyle+ inc. PS, CQG]

![](_page_31_Picture_18.jpeg)

![](_page_31_Figure_19.jpeg)

300

°051

![](_page_31_Figure_20.jpeg)